

Lyman Limit Systems

John M. O'Meara Saint Michael's College

J. Xavier Prochaska, Piero Madau UCSC Hsiao-Wen Chen U. Chicago

October 14, 2010





THANK YOU!





THANK YOU!





THANK YOU!

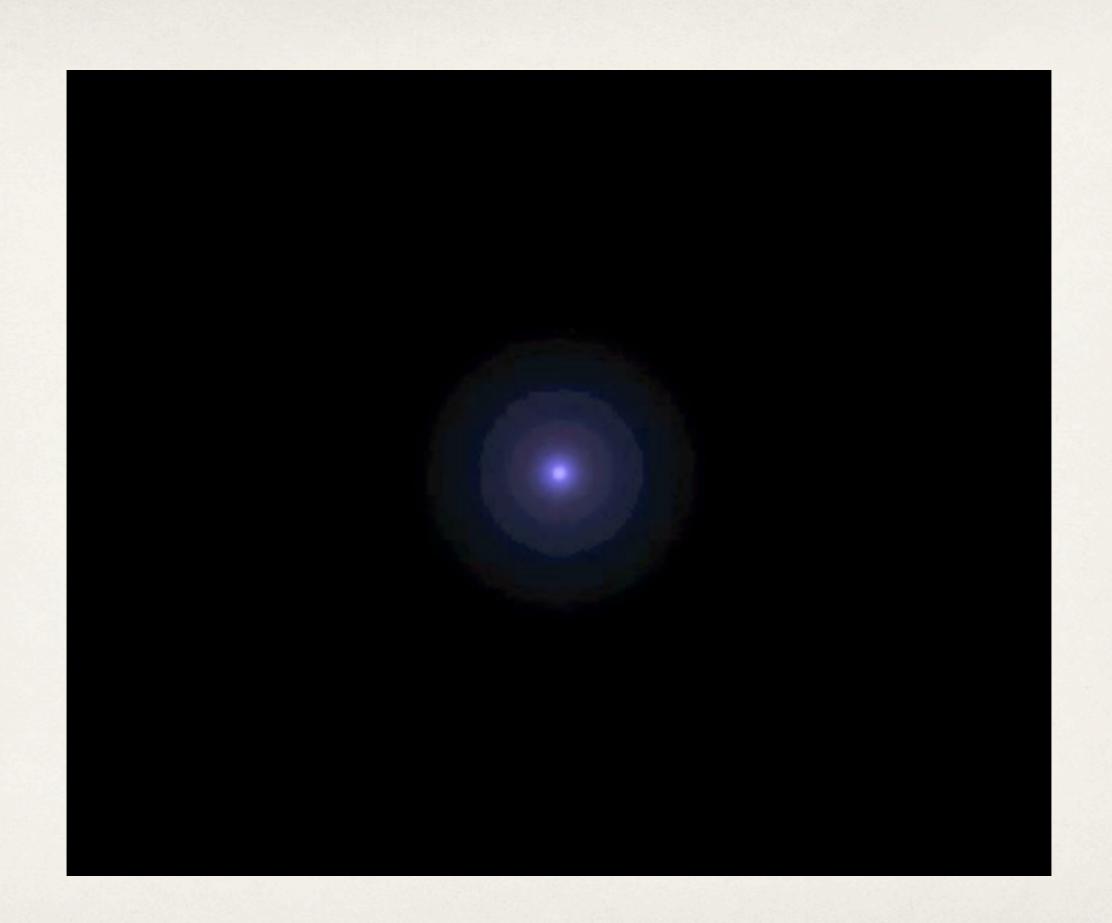


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- * Motivation #3: We still don't know what they are!

How to play the game

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* Quasar surveys easily pick out Lyman limit systems with $\tau_{LL}>2$

A DEFINITIVE SURVEY FOR LYMAN LIMIT SYSTEMS AT $Z\sim 3.5$ WITH THE SLOAN DIGITAL SKY SURVEY

J. XAVIER PROCHASKA¹, JOHN M. O'MEARA², GABOR WORSECK¹ Draft version November 30, 2009

ABSTRACT

We perform a semi-automated survey for $\tau_{912} \ge 2$ Lyman Limit systems (LLSs) in quasar spectra from the Sloan Digital Sky Survey, Data Release 7. From a starting sample of 2473 quasars with $z_{\rm em} = 3.6 - 4.4$, we analyze 469 spectra meeting strict seletion criteria for a total redshift path $\Delta z = 93.8$ and identify 192 intervening systems at $z_{\text{LLS}} \ge 3.3$. The incidence of $\tau_{912} \ge 2$ LLSs per unit redshift, $\ell_{\tau>2}(z)$, is well described by a single-power law at these redshifts: $\ell_{\tau>2}(z) =$ $C_{\text{LLS}}[(1+z)/(1+z_*)]^{\gamma_{\text{LLS}}}$, with $z_* \equiv 3.7$, $C_{\text{LLS}} = 1.9 \pm 0.2$, and $\gamma_{\text{LLS}} = 5.2 \pm 1.5$ (68% c.l.). These values are systematically lower than previous estimates (especially at z < 4) but are consistent with recent measurements of the mean free path to ionizing radiation. Extrapolations of this power-law to z = 0 are inconsistent with previous estimations of $\ell(z)$ at z < 1 and suggest a break at $z \approx 2$, similar to that observed for the Ly α forest. Our results also indicate that the systems giving rise to LLS absorption decrease by $\approx 50\%$ in comoving number density and/or physical size from z = 4to 3.3, perhaps due to an enhanced extragalactic ultraviolet background. The observations place an integral constraint on the H I frequency distribution $f(N_{\rm HI}, X)$ and indicate that the power-law slope $\beta \equiv d \ln f(N_{\rm HI}, X)/d \ln N_{\rm HI}$ is likely shallower than $\beta = -1$ at $N_{\rm HI} \approx 10^{18} \, {\rm cm}^{-2}$. Including other constraints on $f(N_{\rm HI}, X)$ from the literature, we infer that β is steeper than $\beta = -1.7$ at $N_{\rm HI} \approx 10^{15} \,{\rm cm}^{-2}$, implying at least two inflections in $f(N_{\rm HI}, X)$. We also perform a survey for proximate LLSs (PLLSs) and find that $\ell_{\rm PLLS}(z)$ is systematically lower ($\approx 25\%$) than intervening systems. Finally, we estimate that systematic effects impose an uncertainty of 10-20% in the $\ell(z)$ measurements; these effects may limit the precision of all future surveys.

Subject headings: absorption lines – intergalactic medium – Lyman limit systems – SDSS

1. INTRODUCTION

Studies of hydrogen absorption in the lines of sight towards distant quasars have served to both define, and in recent years bring precision to, our cosmological models. The low density, highly ionized Lyman- α forest lines (a.k.a. the intergalactic medium, IGM), with H I column densities $N_{\rm HI} < 10^{17.2} \, {\rm cm}^{-2}$, have through their aggregate statistical properties (e.g. their flux power spectrum, mean flux, and column density distributions) constrained cosmological parameters such as the primordial power spectrum and the baryonic mass density and astrophysical parameters like the amplitude of the ionizing background (e.g. Rauch 1998; Croft et al. 2002; Mc-Donald et al. 2005; Tytler et al. 2004; Faucher-Giguère et al. 2008b). The high-density, predominantly neutral damped Ly α systems (DLAs), with $N_{\rm HI} \ge 10^{20.3} \, {\rm cm}^{-2}$, trace the gas which forms stars, and likely represent the progenitors of modern-day galaxies (e.g. Wolfe et al. 1995, 2005; Prochaska & Wolfe 2009).

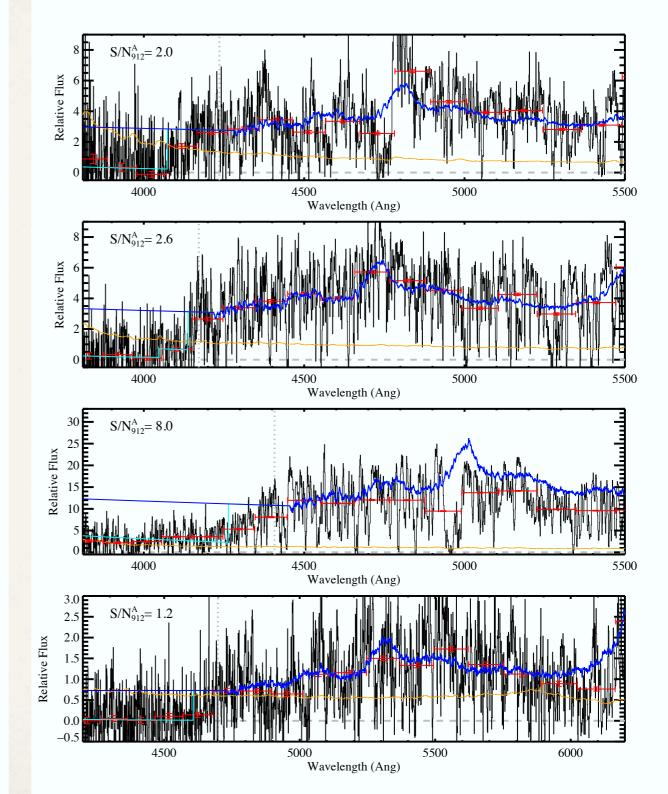
The majority of Lyman– α forest lines and the DLAs have, through analysis of their Ly α lines, precisely measured $N_{\rm HI}$ values that permit detailed study of their physical properties (e.g. metallicity). For systems with intermediate $N_{\rm HI}$ values ($\approx 10^{18} \, {\rm cm}^{-2}$), however, Ly α and most of the Lyman series lines lie on the flat portion of the curve-of-growth making the $N_{\rm HI}$ value difficult to

 1 Department of Astronomy and Astrophysics, UCO/Lick Observatory, University of California, 1156 High Street, Santa Cruz, CA 95064

² Department of Chemistry and Physics, Saint Michael's College One Winooski Park, Colchester, VT 05439 constrain. On the other hand, these systems are optically thick to ionizing radiation and impose a readily identified signature in a quasar spectrum at the Lyman limit. These so-called Lyman limit systems (LLSs), currently the least-well studied of H I absorption systems at high redshift, are the focus of this manuscript.

Historically, the LLSs were among the first class of quasar absorption line (QAL) systems to be surveyed (Tytler 1982). This is because their spectral signature is obvious in low-resolution, low S/N spectra. The principal challenge is that the Lyman limit occurs redward of the atmospheric cutoff only for systems with redshifts z > 2.6. For lower redshifts, one requires spectrometers on space-borne ultraviolet satellites. By the mid 1990's, samples of several tens of LLSs were generated spanning redshifts 0 < z < 4 (Sargent et al. 1989; Lanzetta 1991; Storrie-Lombardi et al. 1994; Stengler-Larrea et al. 1995). These results were derived from heterogeneous sets of quasars discovered from a combination of color-selection, radio detection, and slitless spectroscopic surveys. The spectra, too, were acquired with a diverse set of instrumentation and therefore varying S/N and spectral resolution mitigating differing sensitivity to the precise optical depth at the Lyman limit. Although the results were not fully consistent with one another, the general picture that resulted was a rapidly evolving population of absorption systems reasonably described by a $(1+z)^{1.5}$ power-law.

Cosmologically, the LLSs contribute much if not most of the universe's opacity to ionizing radiation. And, until recently, the observed incidence of the LLS provided the only direct means of estimating the mean free path $\lambda_{\rm mfp}^{912}$



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THE SDSS DAMPED Ly α SURVEY: DATA RELEASE 3

JASON X. PROCHASKA & STÉPHANE HERBERT-FORT Department of Astronomy and Astrophysics, UCO/Lick Observatory; University of California, 1156 High Street, Santa Cruz, CA 95064; xavier@ucolick.org, shf@ucolick.org AND

ARTHUR M. WOLFE Department of Physics, and Center for Astrophysics and Space Sciences, University of California, San Diego, Gilman Dr., La Jolla; CA 92093-0424; awolfe@ucsd.edu Accepted to ApJ: August 15, 2005

ABSTRACT

We present the results from a damped $Ly\alpha$ survey of the Sloan Digital Sky Survey, Data Release 3. We have discovered over 500 new damped Ly α systems at z > 2.2 and the complete statistical sample for z > 1.6 has more than 600 damped Ly α galaxies. We measure the H I column density distribution $f_{\rm HI}(N, X)$ and its zeroth and first moments (the incidence $\ell_{\rm DLA}$ and gas mass-density $\Omega_a^{\rm DLA}$ of damped $Ly\alpha$ systems, respectively) as a function of redshift. The key results include: (1) the full SDSS-DR3 $f_{\rm HI}(N, X)$ distribution ($z \sim 3.06$) is well fit by a Γ -function (or double power-law) with 'break' column density $N_{\gamma} = 10^{21.5\pm0.1} \text{ cm}^{-2}$ and 'faint-end' slope $\alpha = -1.8\pm0.1$; (2) the shape of the $f_{\text{HI}}(N,X)$ distributions in a series of redshift bins does not show evolution; (3) the incidence and gas mass density of damped systems decrease by $35 \pm 9\%$ and $50 \pm 10\%$ during $\approx 1 \,\text{Gyr}$ between the redshift interval z = [3., 3.5] to z = [2.2, 2.5]; and (4) the incidence and gas mass density of damped Ly α systems in the lowest SDSS redshift bin (z = 2.2) are consistent with the current values. We investigate a number of systematic errors in damped Ly α analysis and identify only one important effect: we measure $40 \pm 20\%$ higher Ω_a^{DLA} values toward a subset of brighter quasars than toward a faint subset. This effect is contrary to the bias associated with dust obscuration and suggests that gravitational lensing may be important. Comparing the results against several models of galaxy formation in Λ CDM, we find all of the models significantly underpredict ℓ_{DLA} at z = 3 and only SPH models with significant feedback (Nagamine et al.) may reproduce Ω_g^{DLA} at high redshift. Based on our results for the damped Ly α systems, we argue that the Lyman limit systems contribute $\approx 33\%$ of the universe's H I atoms at all redshifts z = 2 to 5. Furthermore, we infer that the $f_{\rm HI}(N, X)$ distribution for $N_{\rm HI} < 10^{20} \,{\rm cm}^{-2}$ has an inflection with slope $d\log f/d\log N > -1$. We advocate a new mass density definition – the mass density of predominantly neutral gas Ω_q^{Neut} – to be contrasted with the mass density of gas associated with H I atoms. We contend the damped Ly α systems contribute > 80% of Ω_a^{Neut} at all redshifts and therefore are the main reservoirs for star formation.

${\it Subject\ headings:\ galaxies:\ evolution\ --\ intergalactic\ medium\ --\ quasars:\ absorption\ lines}$

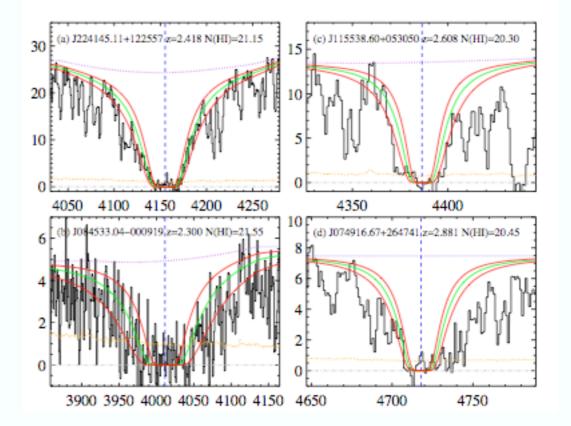
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1. INTRODUCTION

The damped Ly α systems are the class of quasar absorption line systems with H I column density $N_{\rm HI} \geq 2 \times 10^{20} \, {\rm cm}^{-2}$ (see Wolfe, Gawiser, & Prochaska 2005, for a review). Unlike the Ly α forest, the damped Ly α systems are comprised of predominantly neutral gas and are proposed to be the progenitors of galaxies like the Milky Way (e.g. Kauffmann 1996). Wolfe et al. (1986) established the $N_{th} = 2 \times 10^{20} \, {\rm cm}^{-2}$ threshold primarily to correspond to the surface density limit of local 21 cm observations at that time. It is somewhat fortuitous that this threshold roughly corresponds to the transition from primarily ionized gas to predominantly neutral gas (e.g. Viegas 1995; Prochaska & Wolfe 1996; Prochaska 1999; Vladilo et al. 2001).

For the past two decades, several groups have surveyed high z quasars for the damped Ly α systems (Wolfe et al. 1986, 1995; Storrie-Lombardi et al. 1996a; Storrie-Lombardi & Wolfe 2000; Péroux et al. 2003). These surveys measured the H I frequency distribution function and its moments: the incidence of the damped Ly α systems $\ell_{\rm DLA}$ and the gas mass density of these galaxies $\Omega_g^{\rm DLA}$. The latter quantity has cosmological significance. Its evolution constrains the build-up of structure within hierarchical cosmology (e.g. Ma & Bertschinger 1994; Klypin et al. 1995), it serves as the important neutral gas reservoir for star formation at high redshift and describes the competition between gas accretion and star formation (e.g. Fall & Pei 1993), and it constrains models of galaxy formation in hierarchical cosmology (e.g. Somerville, Primack, & Faber 2001; Cen et al. 2003; Nagamine, Springel, & Hernquist 2004). Previous surveys have reported statistical error on Ω_g^{DLA} of $\approx 30\%$ in redshift intervals $\Delta z \approx 0.5$ at high redshift. As we enter the so-called era of precision cosmology, we aspire to constrain Ω_g^{DLA} to better than 10%. Although not formally a cosmological parameter, a precise determination of Ω_g^{DLA} and its redshift evolution are fundamental constraints on any cosmological theory of galaxy formation.

In Prochaska & Herbert-Fort (2004), hereafter PH04, we initiated a survey for the damped Ly α systems in the quasar spectra of the Sloan Digital Sky Survey (SDSS) Data Release 1 (DR1). We demonstrated that the spectral resolution, signal-to-noise ratio (SNR), and wavelength coverage of the SDSS spectra are well suited to survey the damped Ly α systems at z > 2.2. We reported on the number of damped Ly α systems per unit redshift and the



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- * The Lyman Alpha forest is well sampled

Difficulties

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Difficulties

* Going below z~2.7 means going to space

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- * Going below z~2.7 means going to space
- * The integrated opacity includes (and is probably dominated by) $\tau_{LL} \sim 1$

The good news

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The good news

- * Going below z~2.7 means going to space!
- * The universe expands!

Build a homogenous sample of QSOs from SDSS

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- Bright, z~2, no BAL

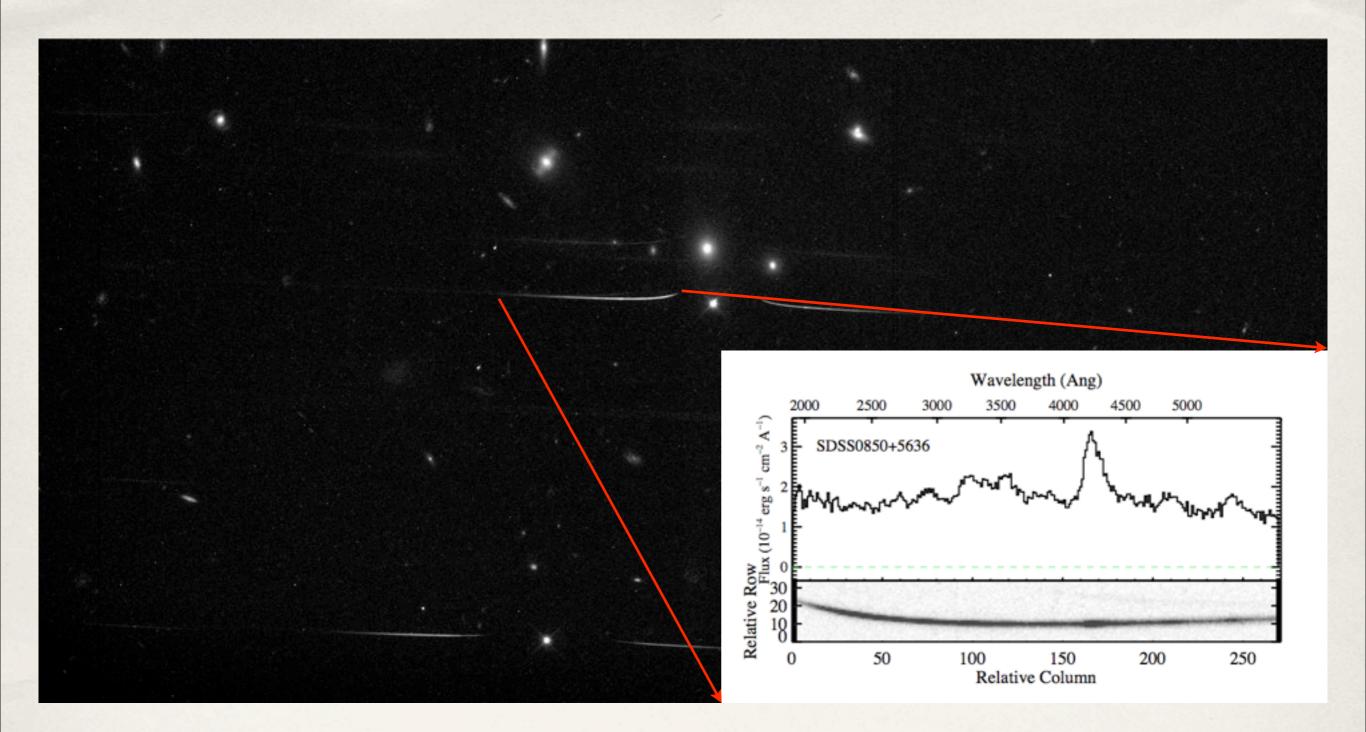
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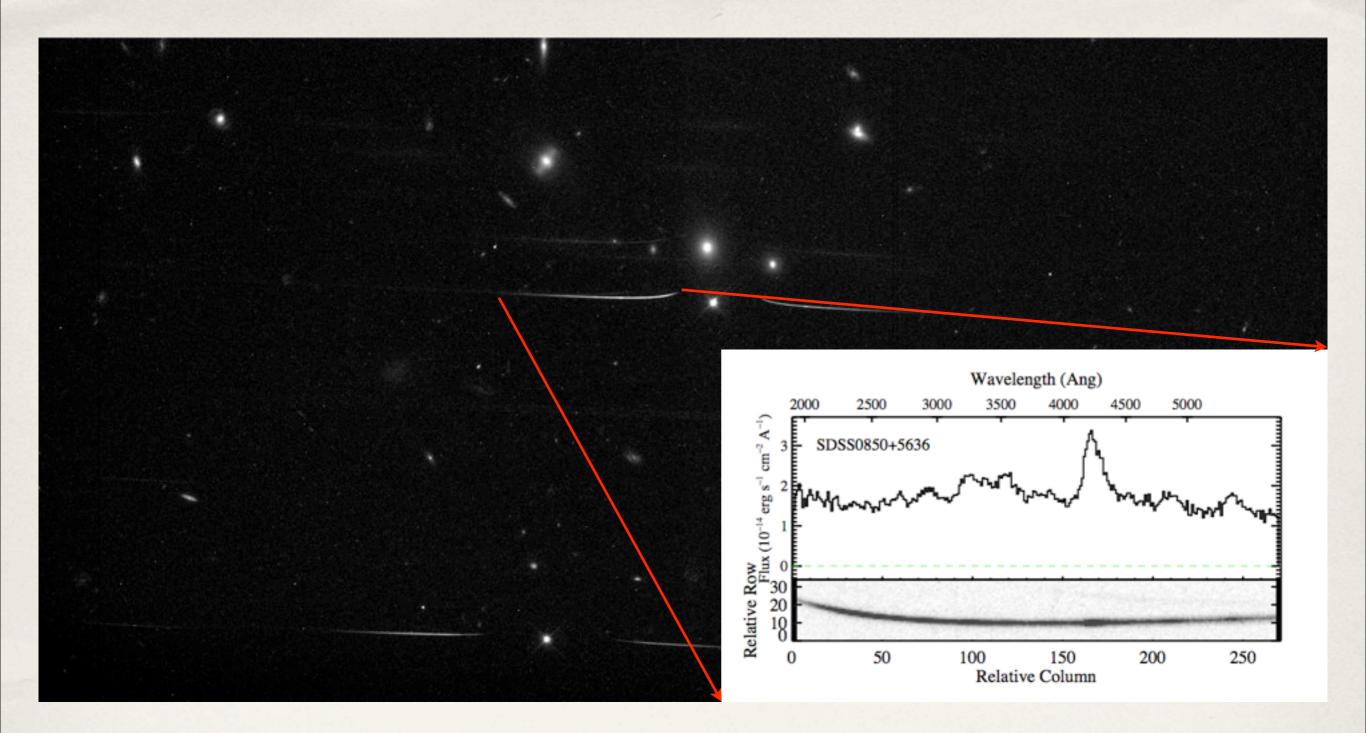
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- * Redshift path at *z*~2 with WFC3 as big as SDSS at *z*~4!



SDSS0850+5636

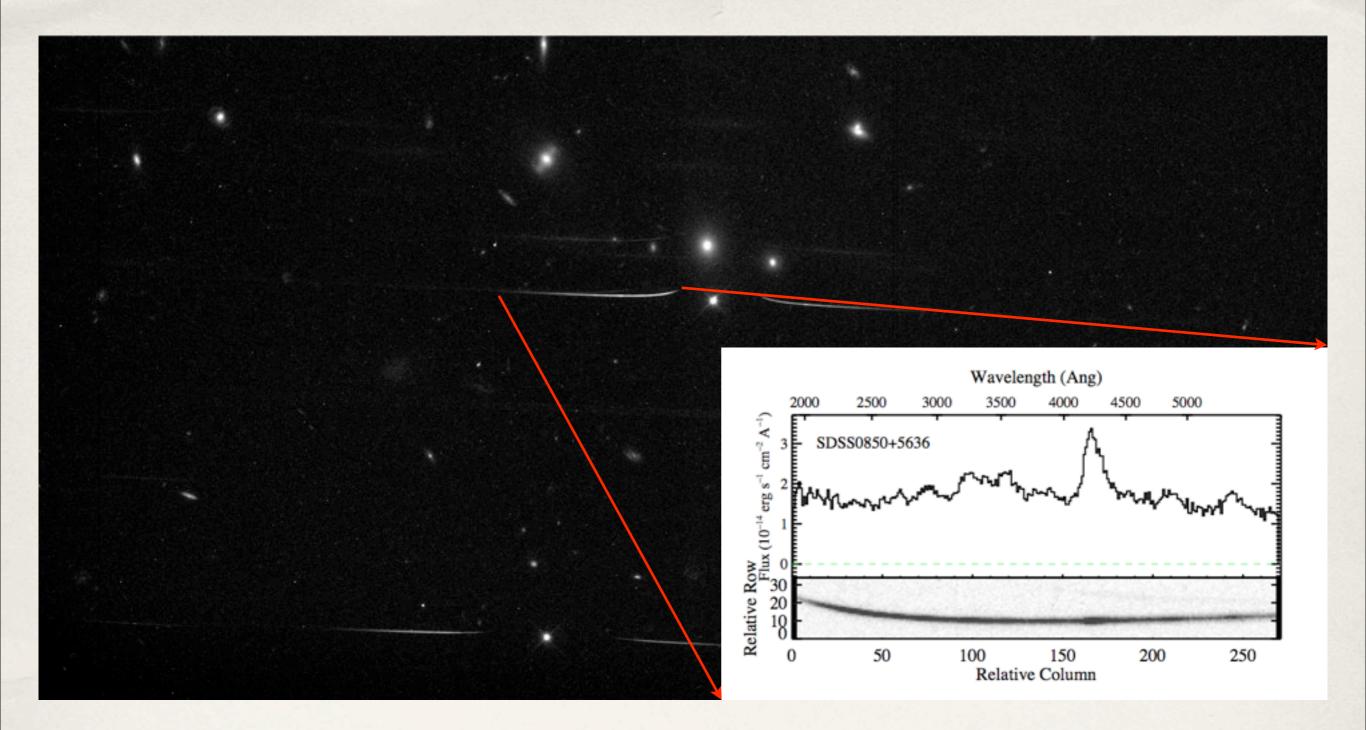


SDSS0850+5636



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Thanks ST-ECF folks!



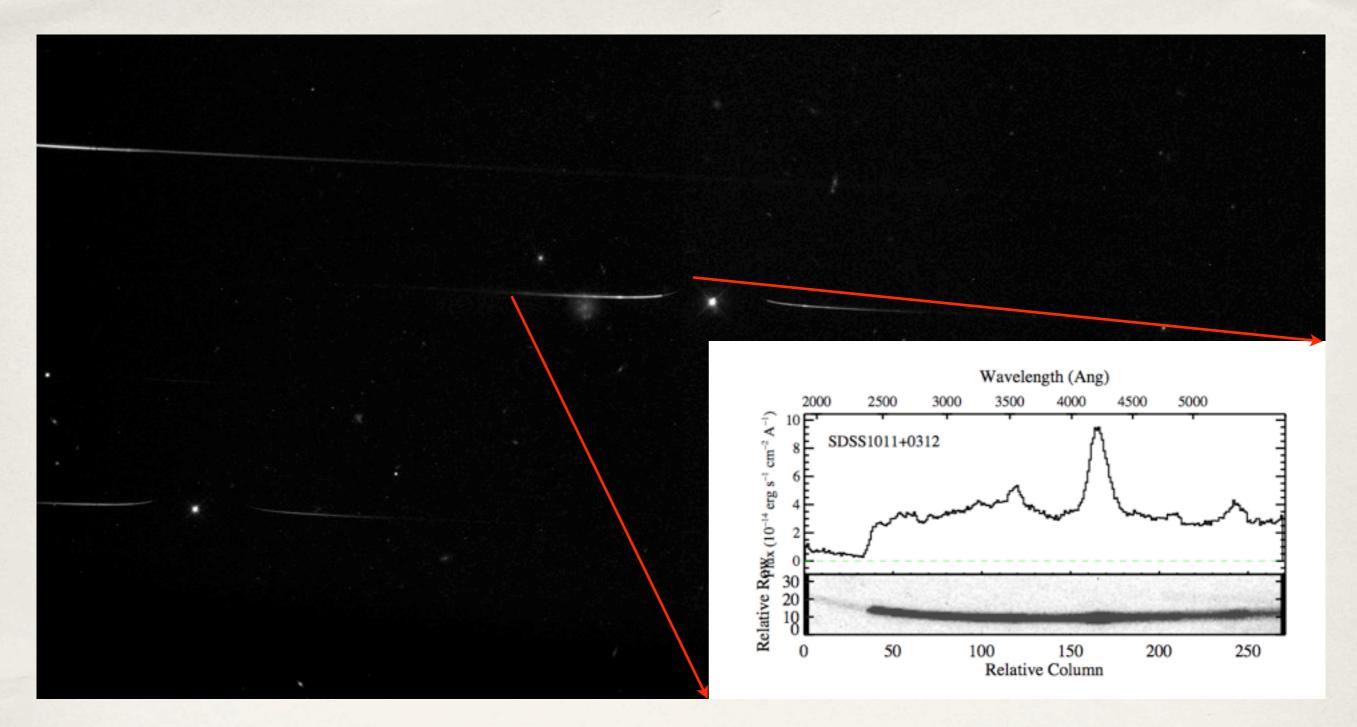
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Thanks ST-ECF folks!

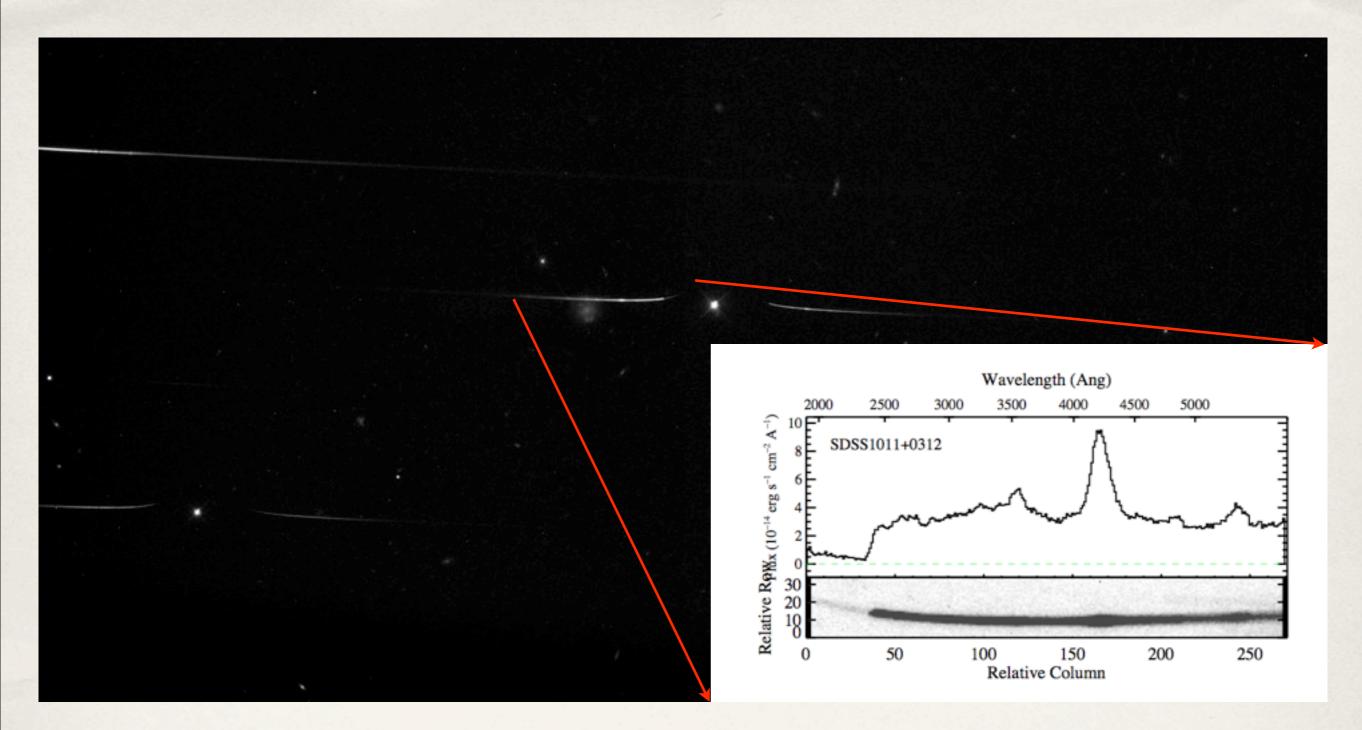
No LLS



SDSS1011 + 0312

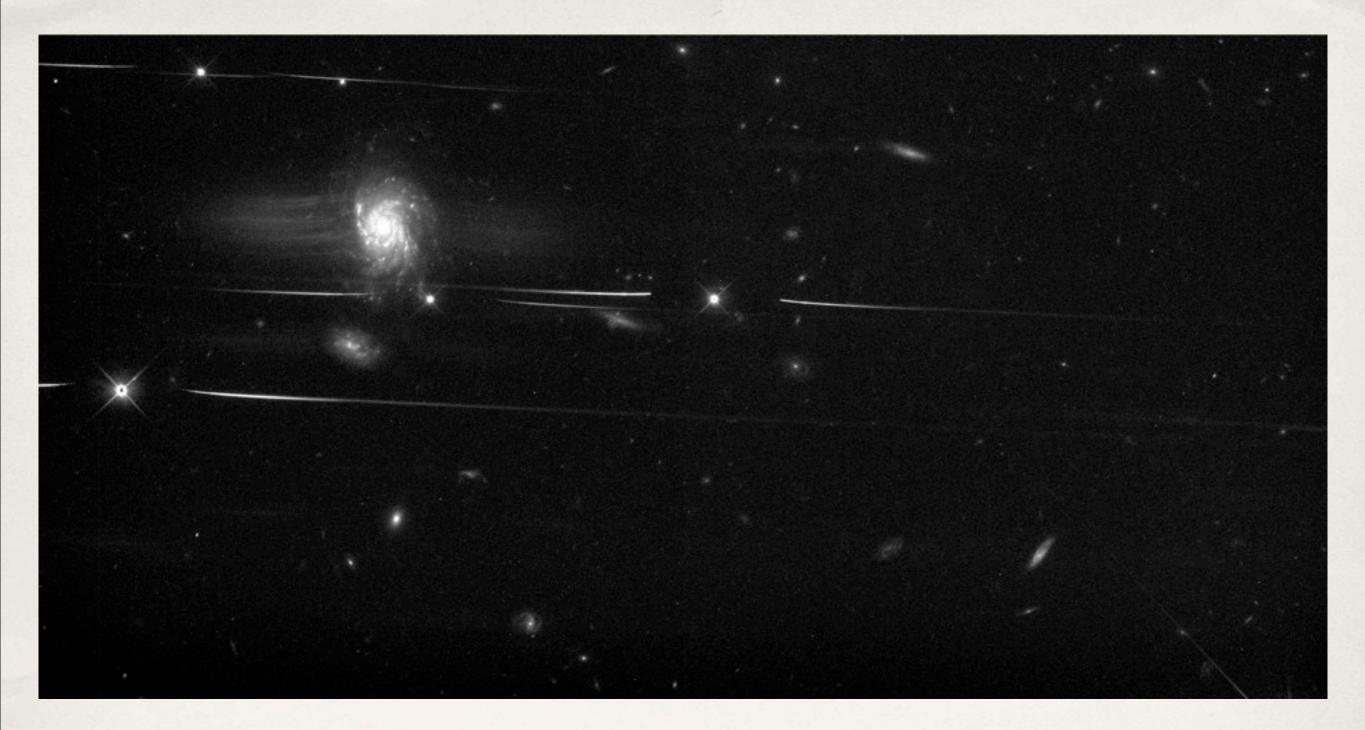


SDSS1011 + 0312

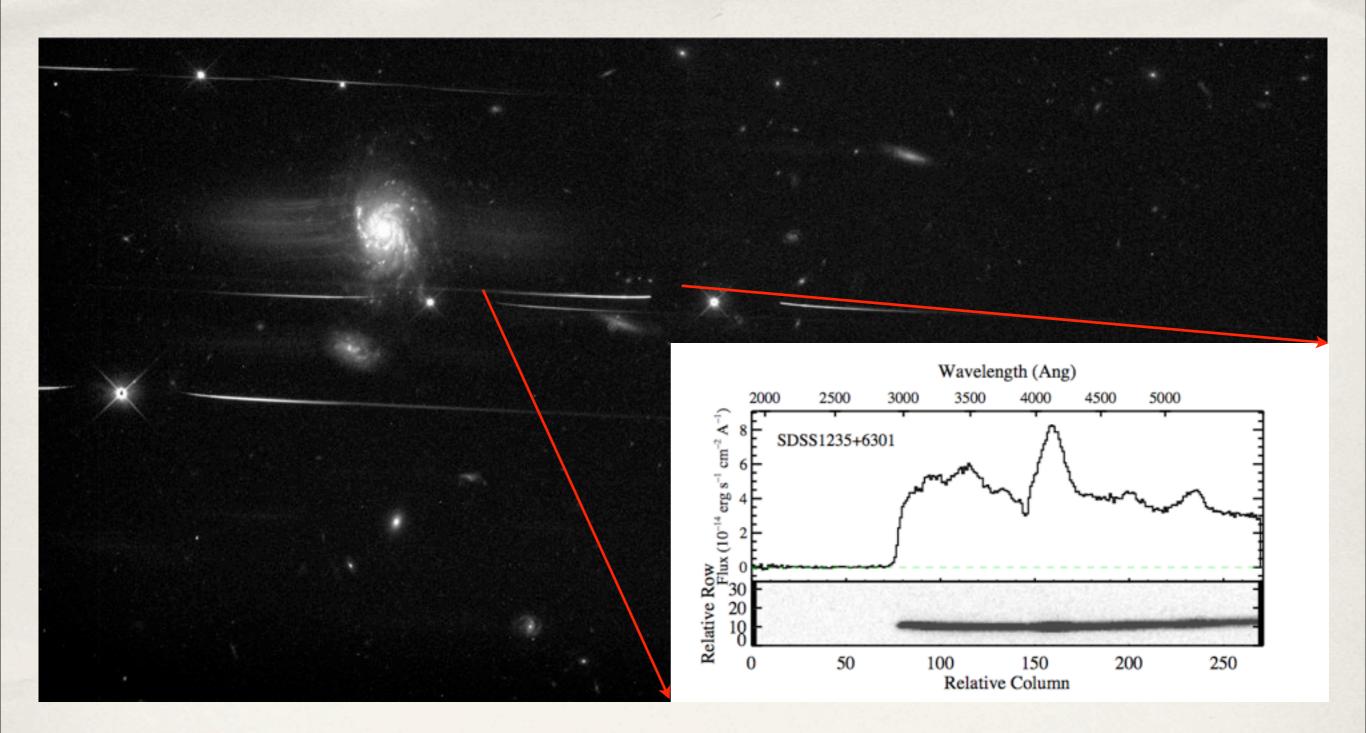


SDSS1011 + 0312

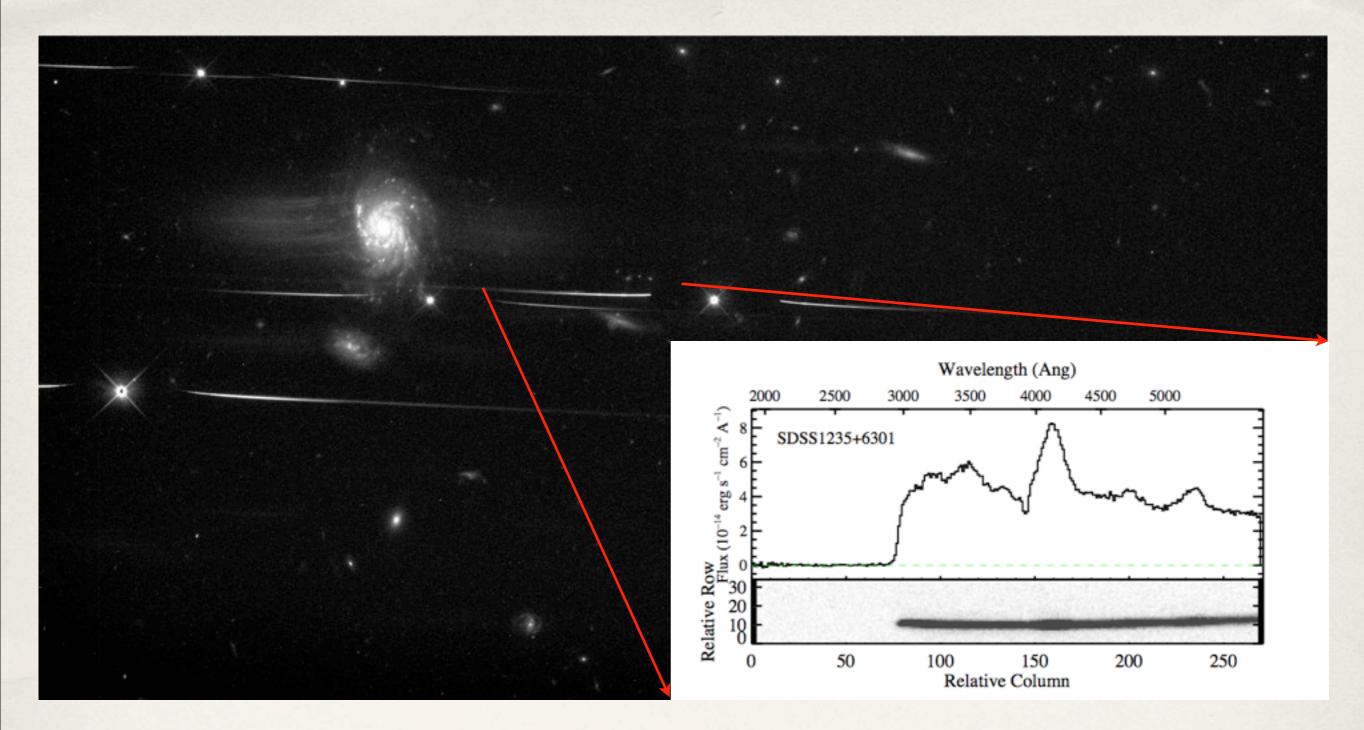
Partial LLS



SDSS1235 + 6301



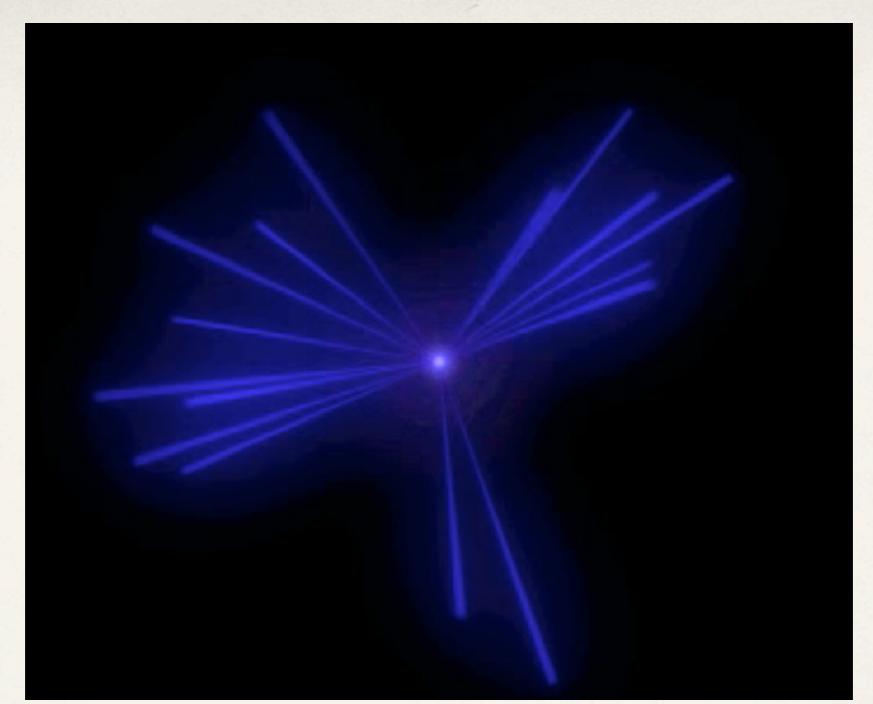
SDSS1235+6301



SDSS1235+6301

Full LLS

IGM Opacity-->The Mean Free Path



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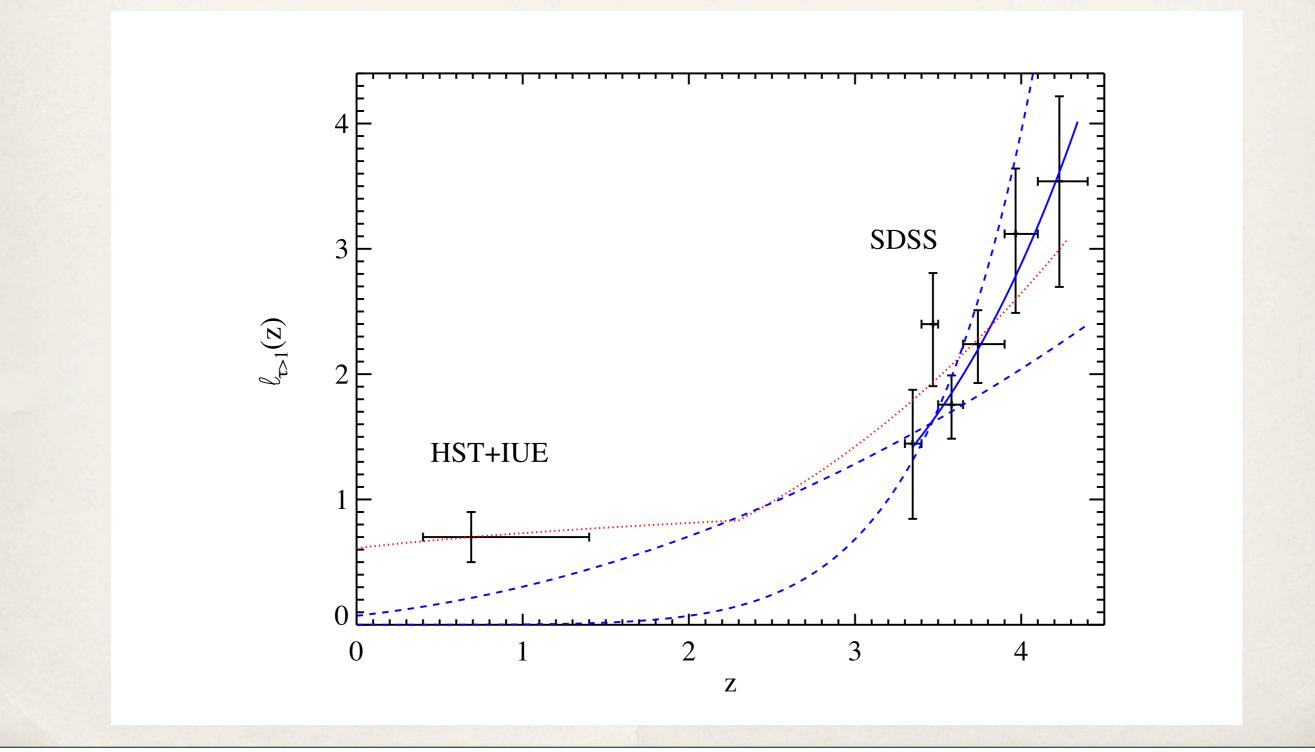
Begin with counting the LLS

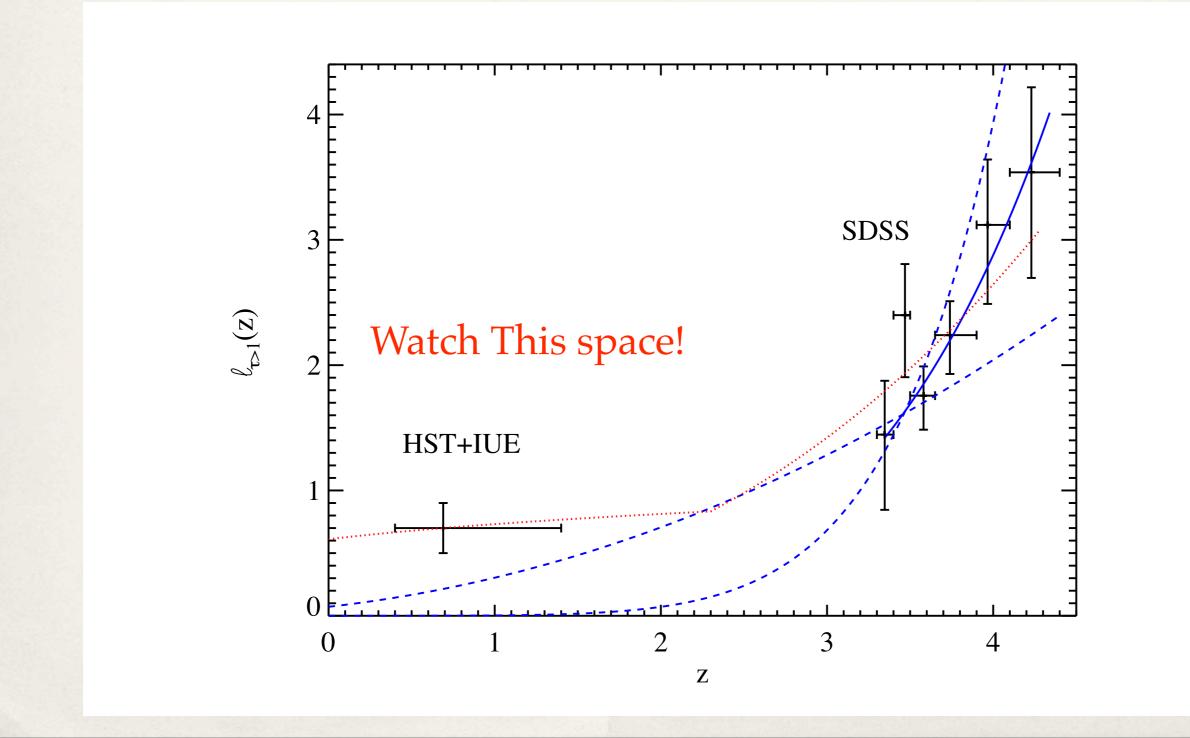
- * Begin with counting the LLS
- * Determine the N(HI) frequency distribution

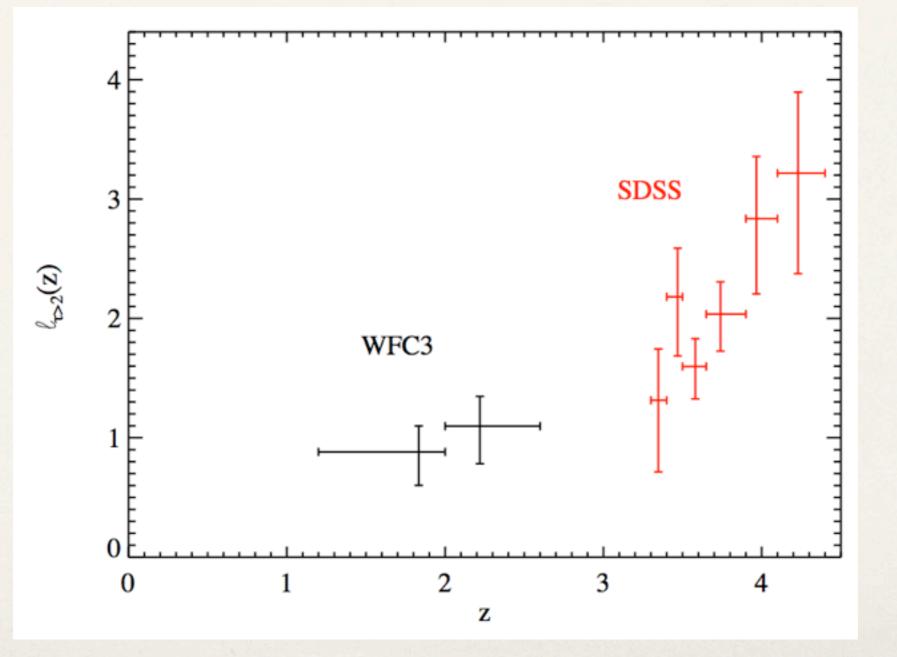
- * Begin with counting the LLS
- Determine the N(HI) frequency distribution
- * Integrate them up!

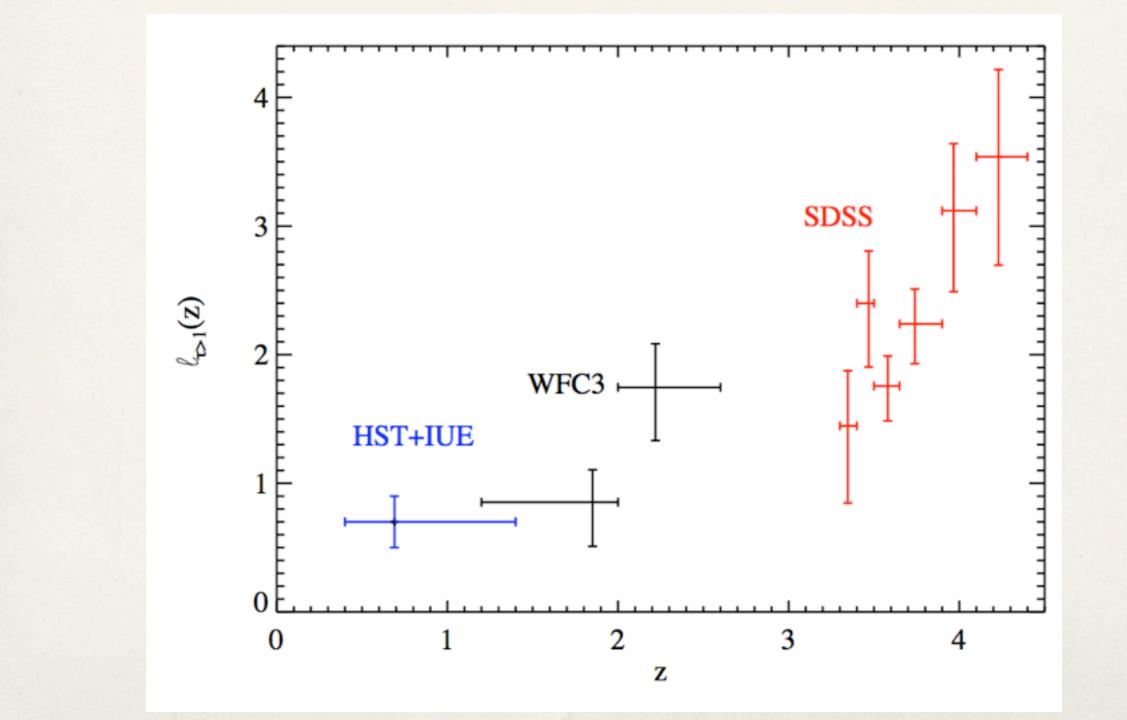
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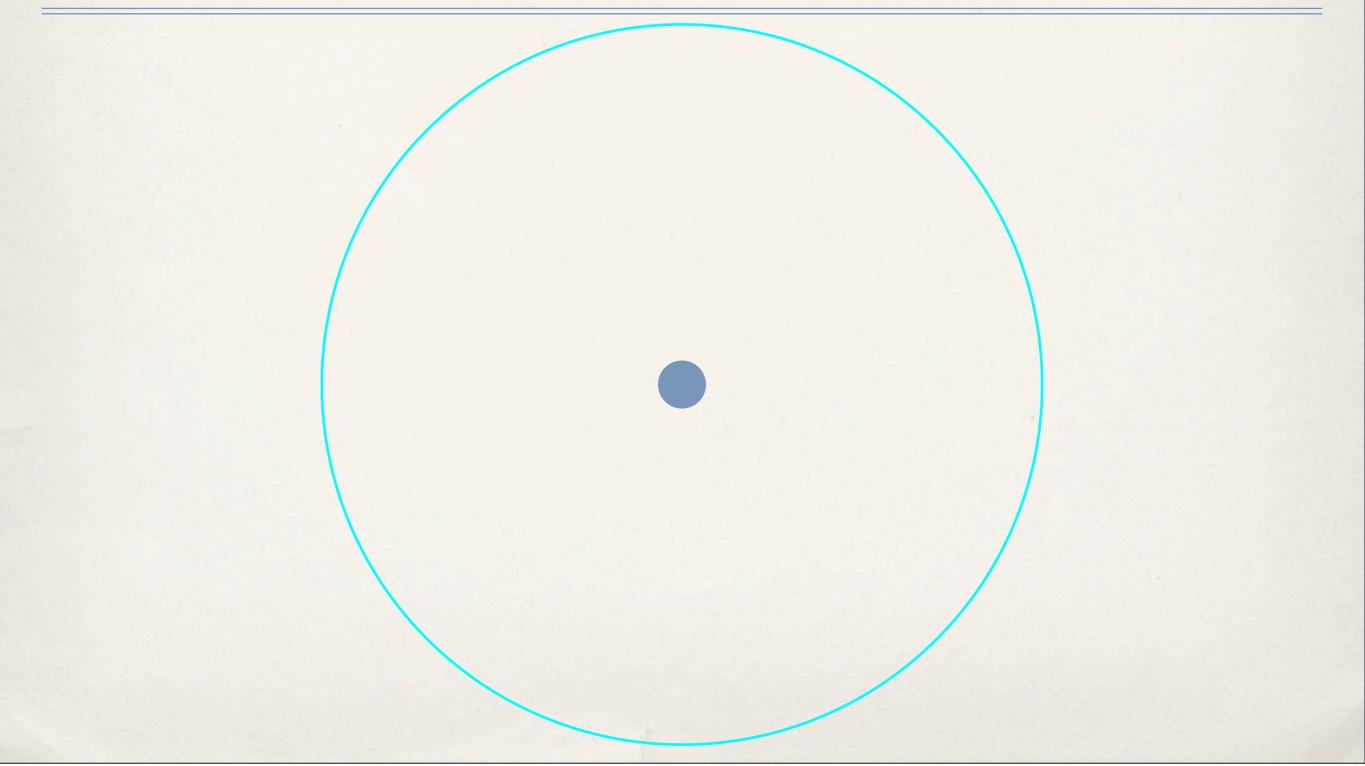
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m eff,LL}(z_{912},z_q) = \int\limits_{z_{912}}^{z_q} \int\limits_{0}^{\infty} f(N_{
m HI},z') \{1 - \exp\left[-N_{
m HI}\sigma_{
m ph}(z')
ight] \} dN_{
m HI} dz'$$

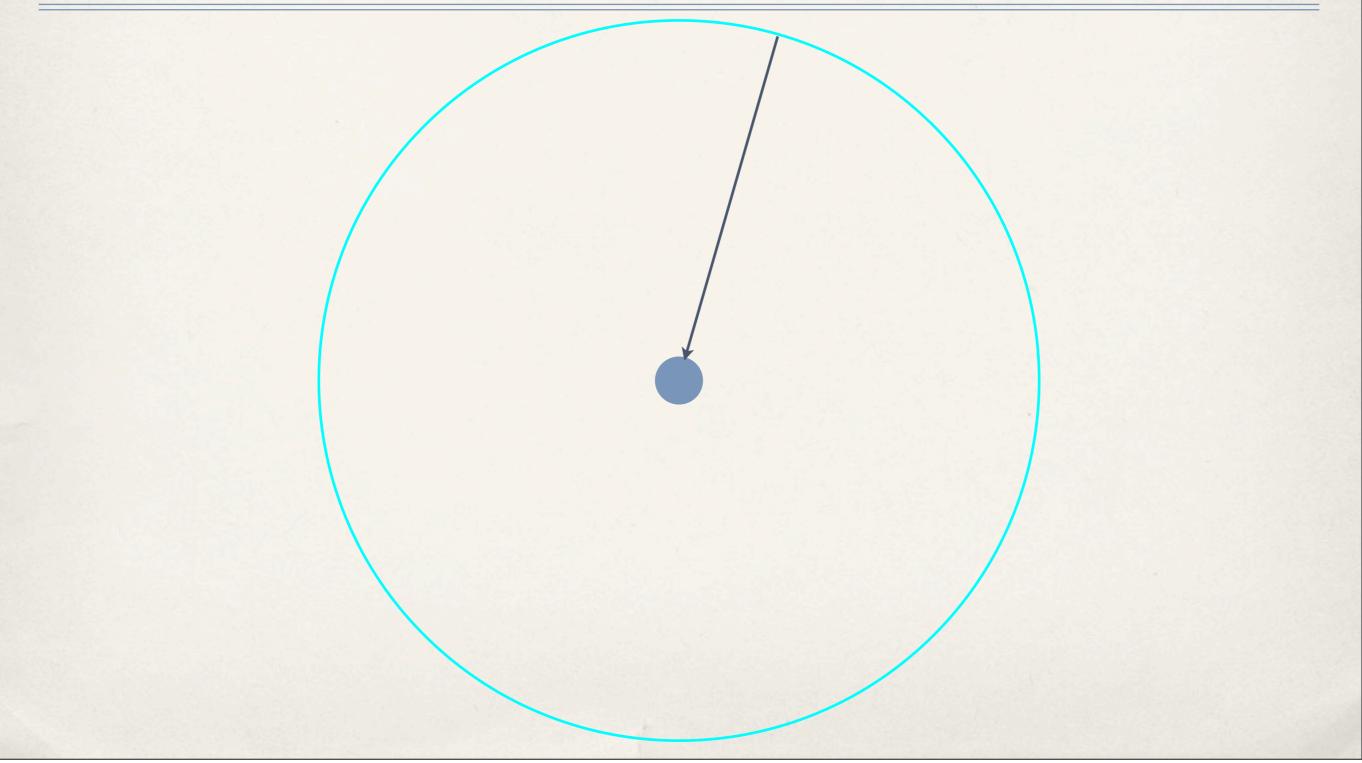


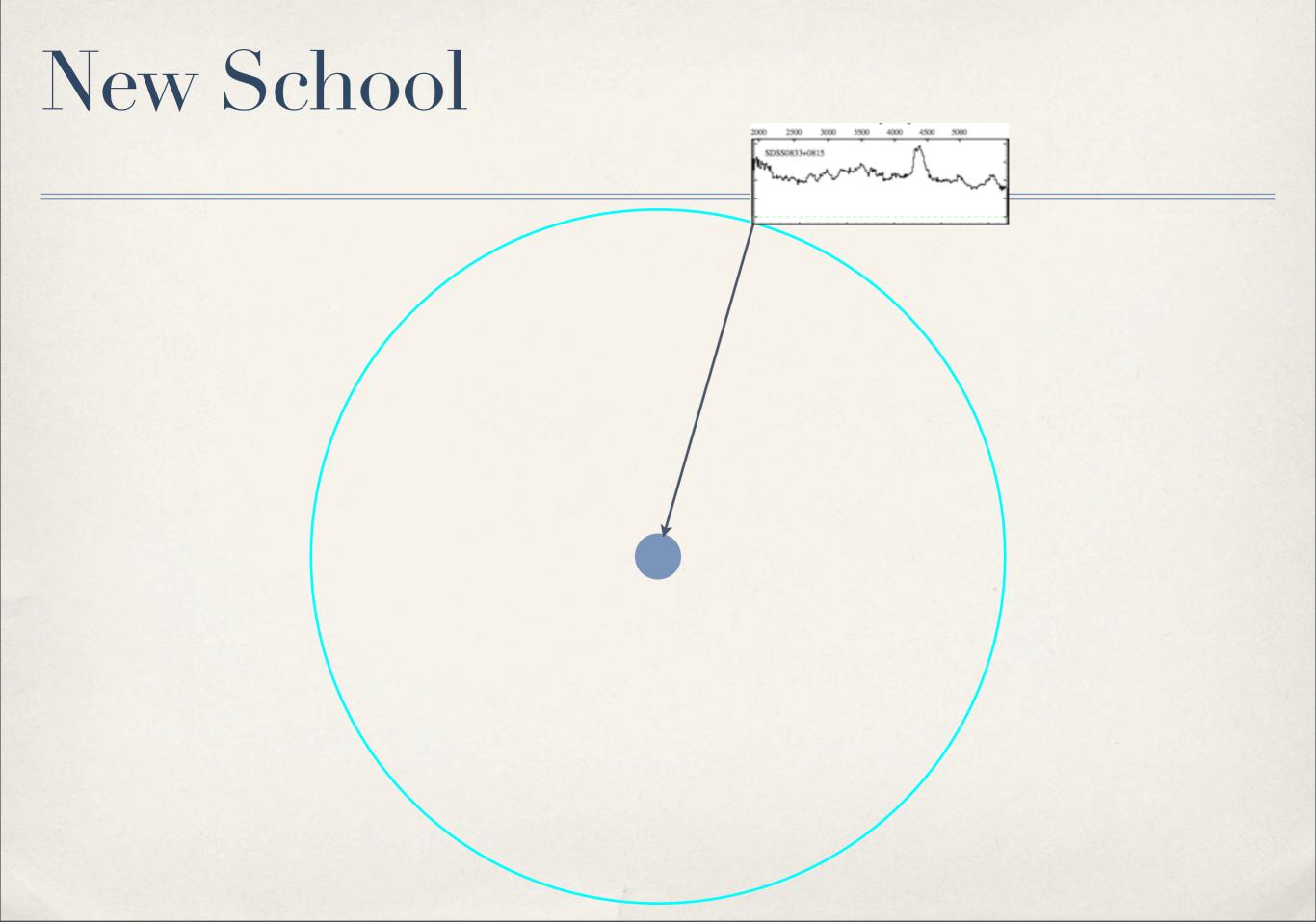


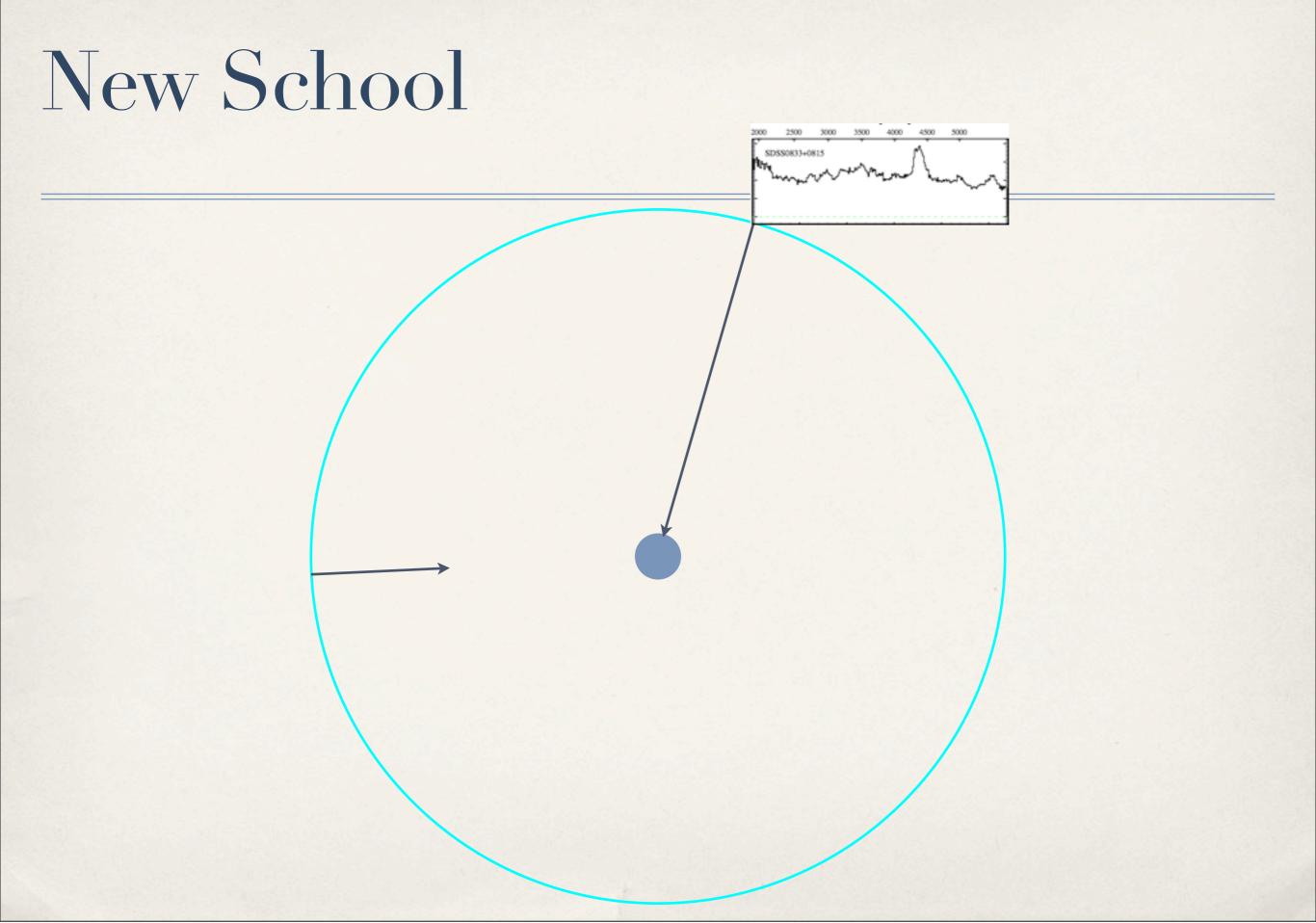


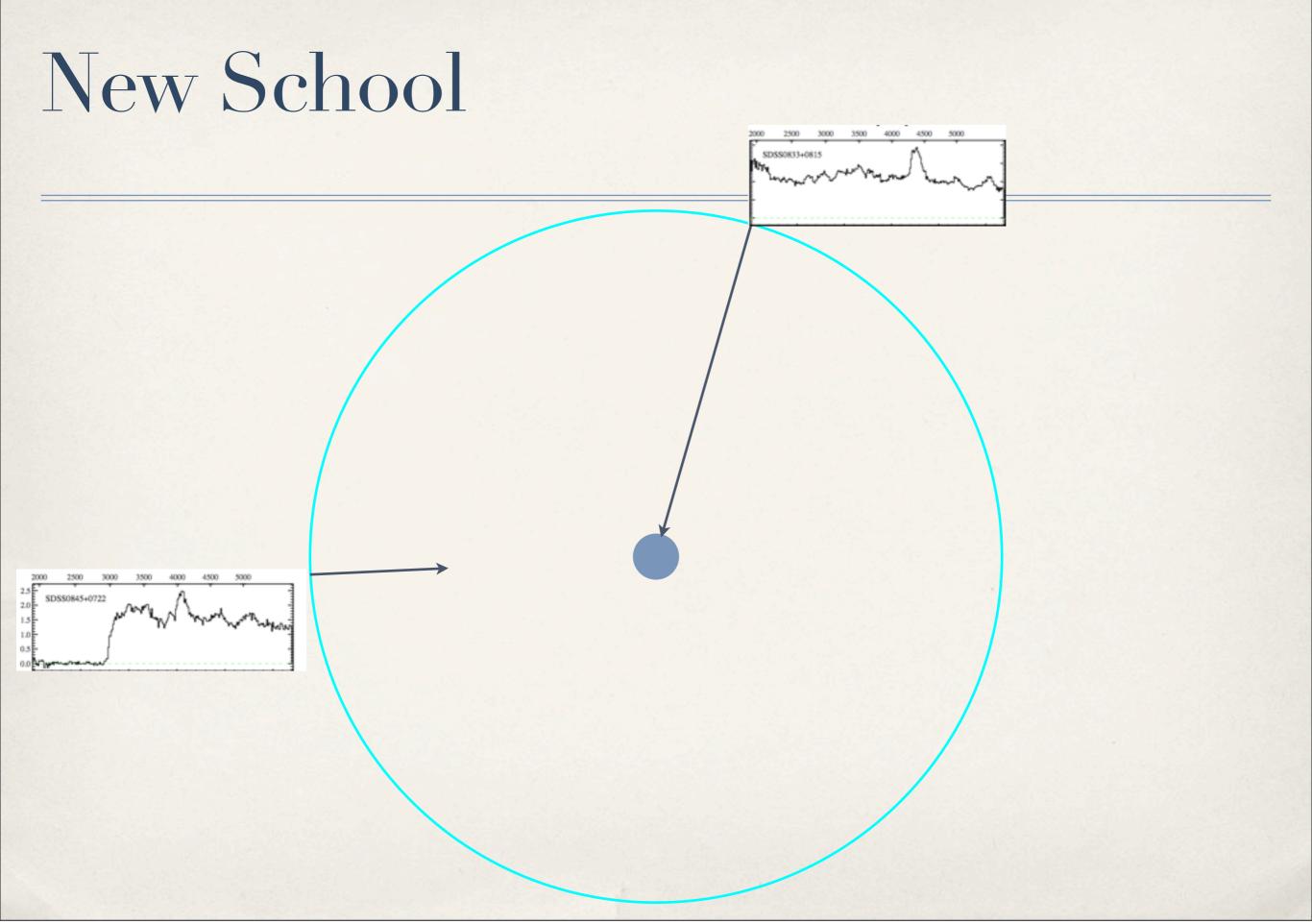


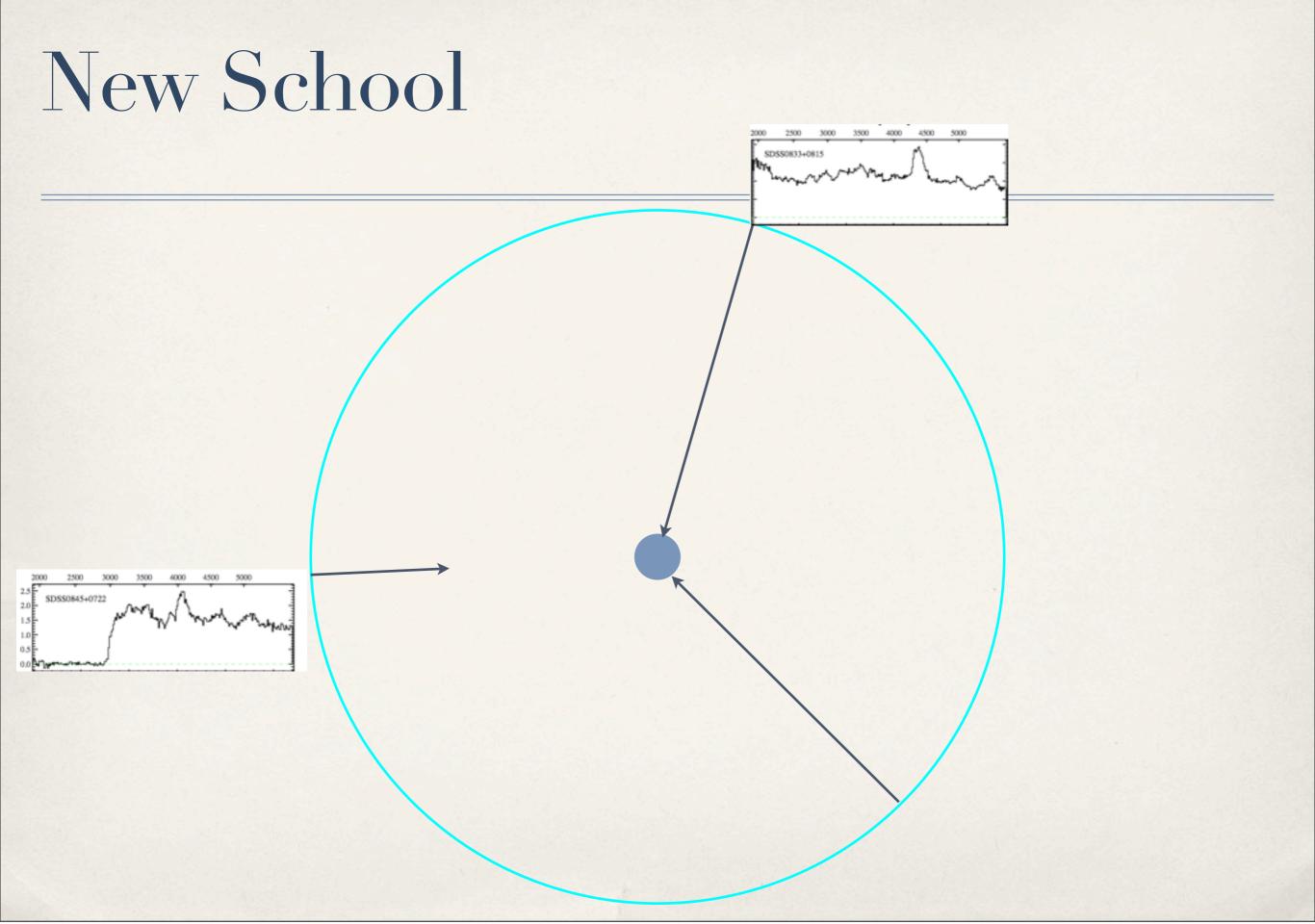


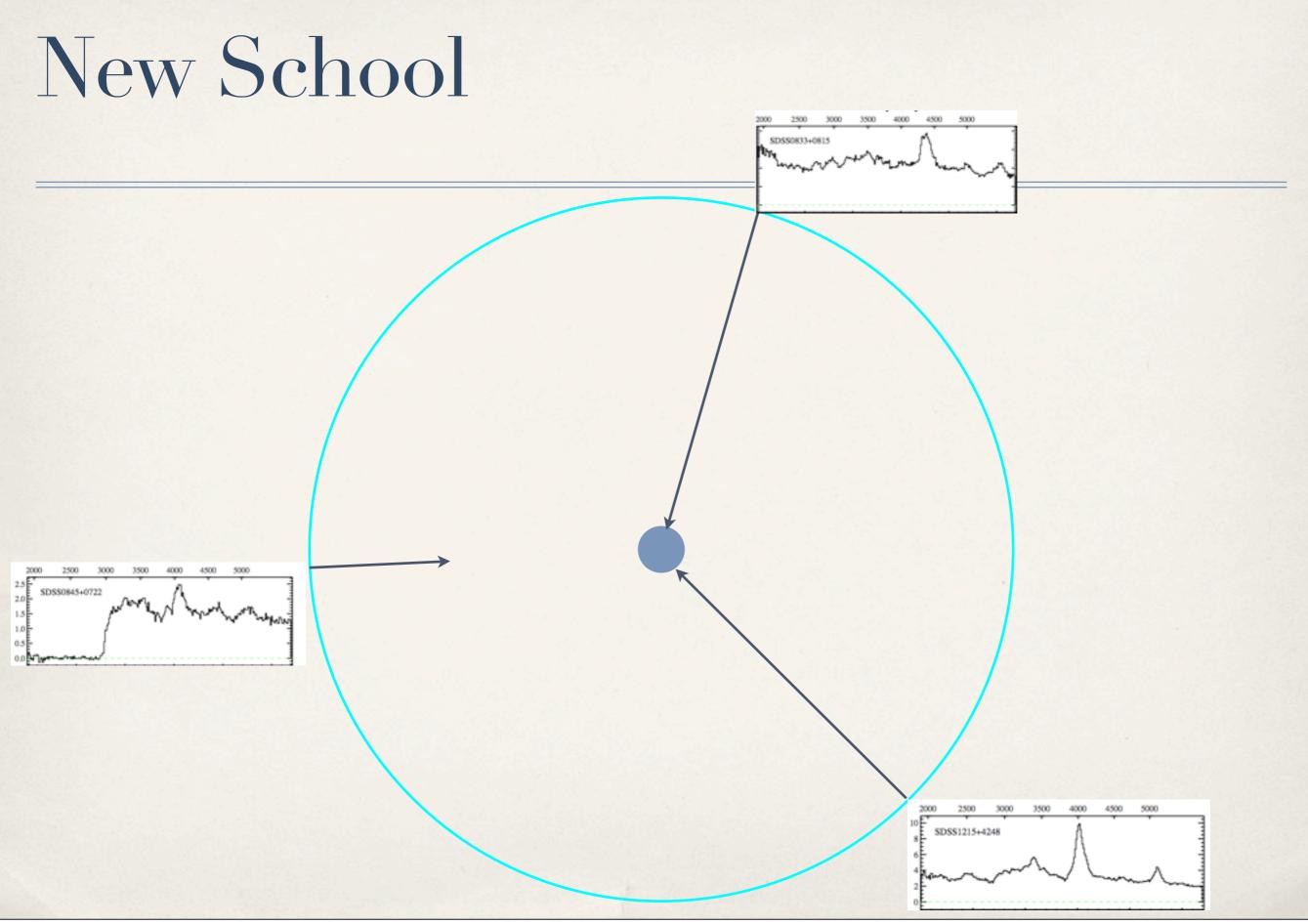


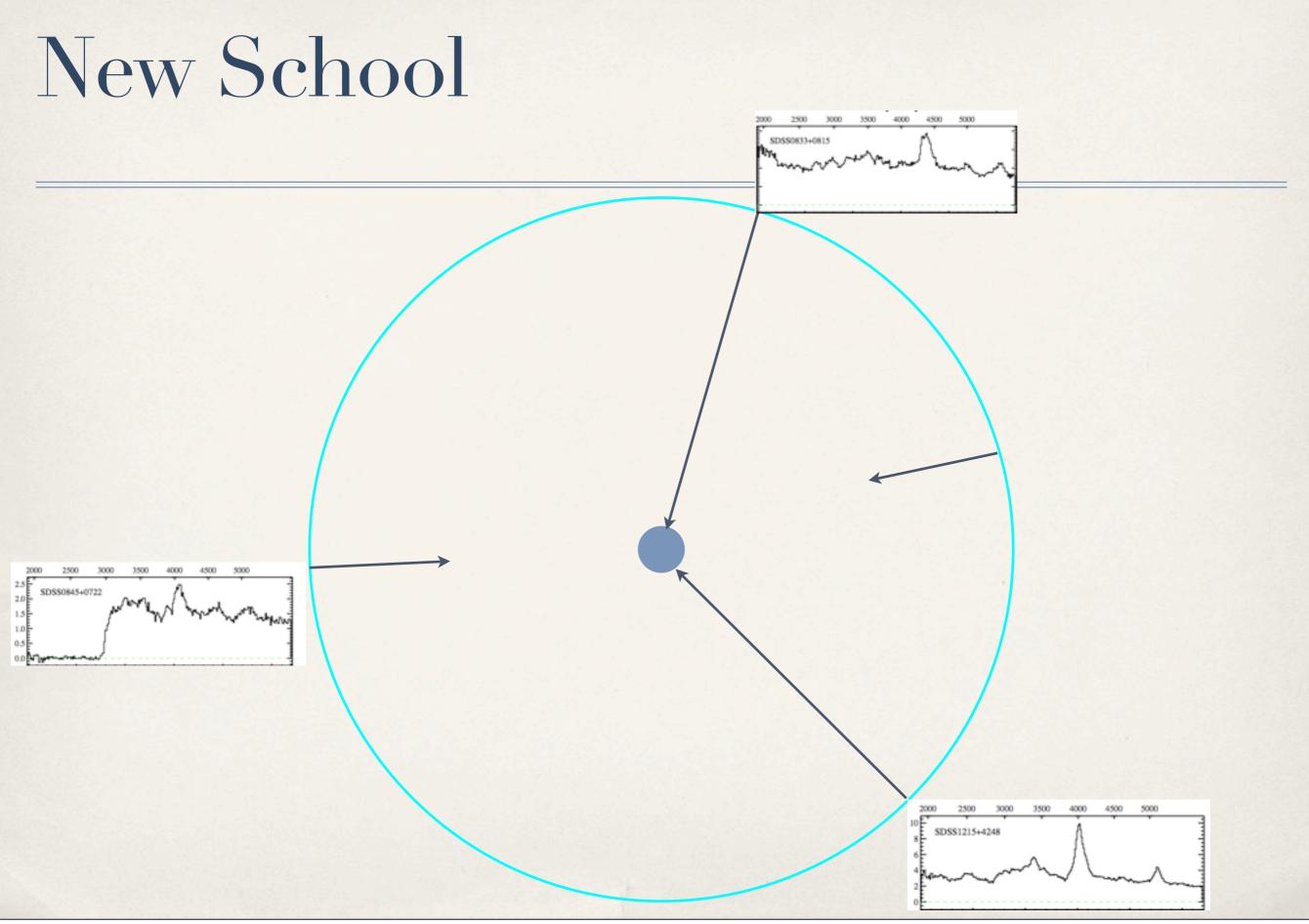


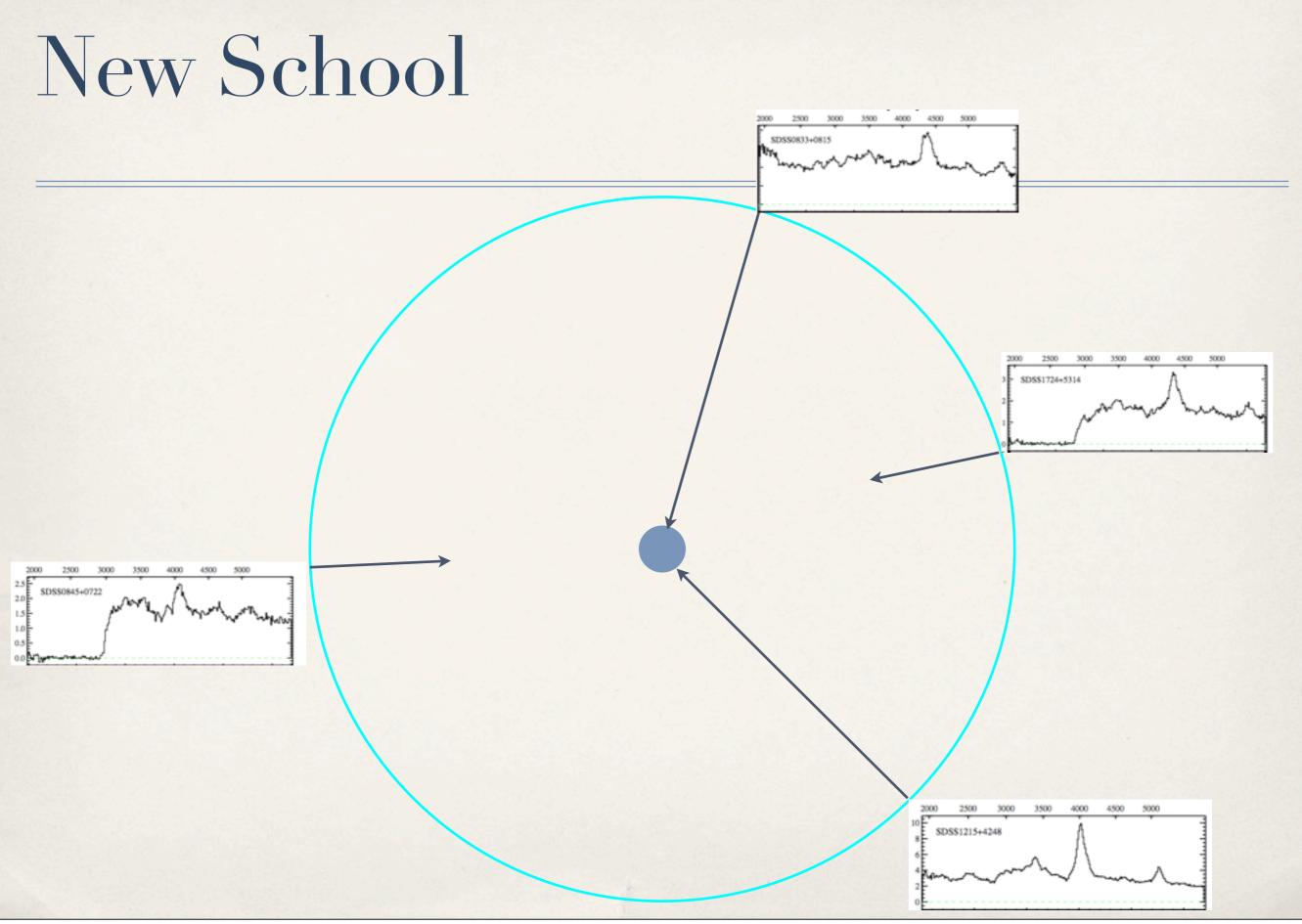


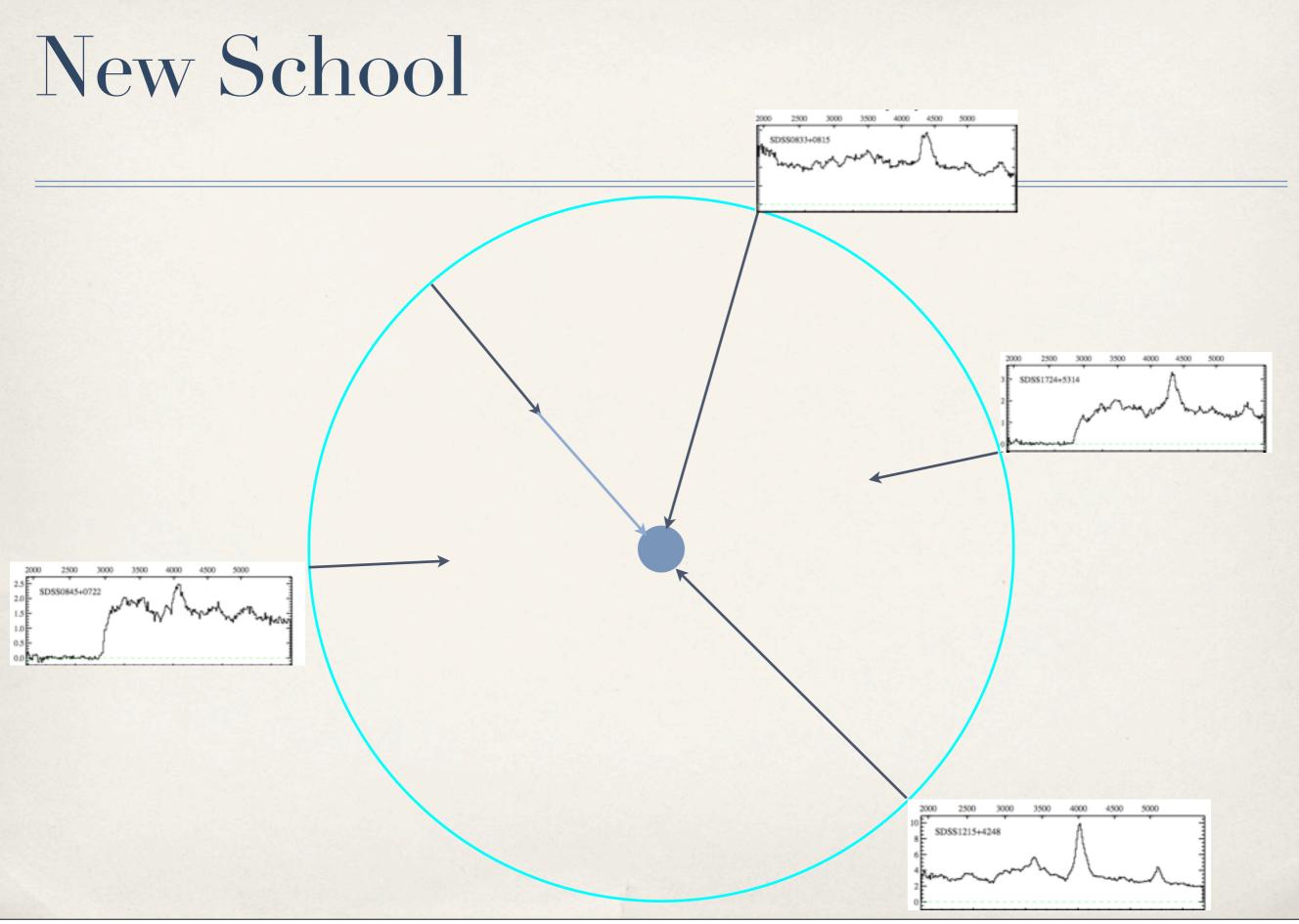


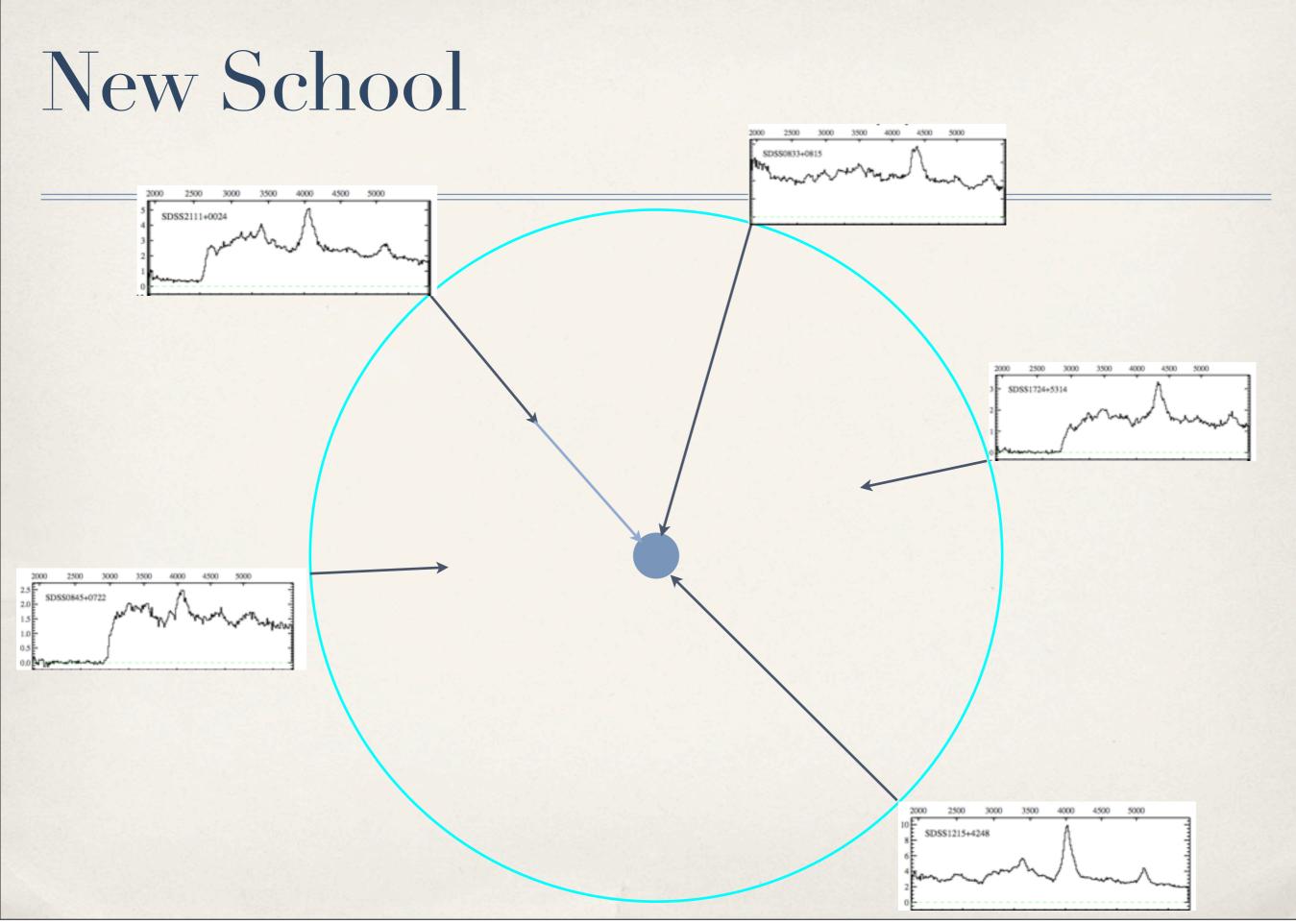


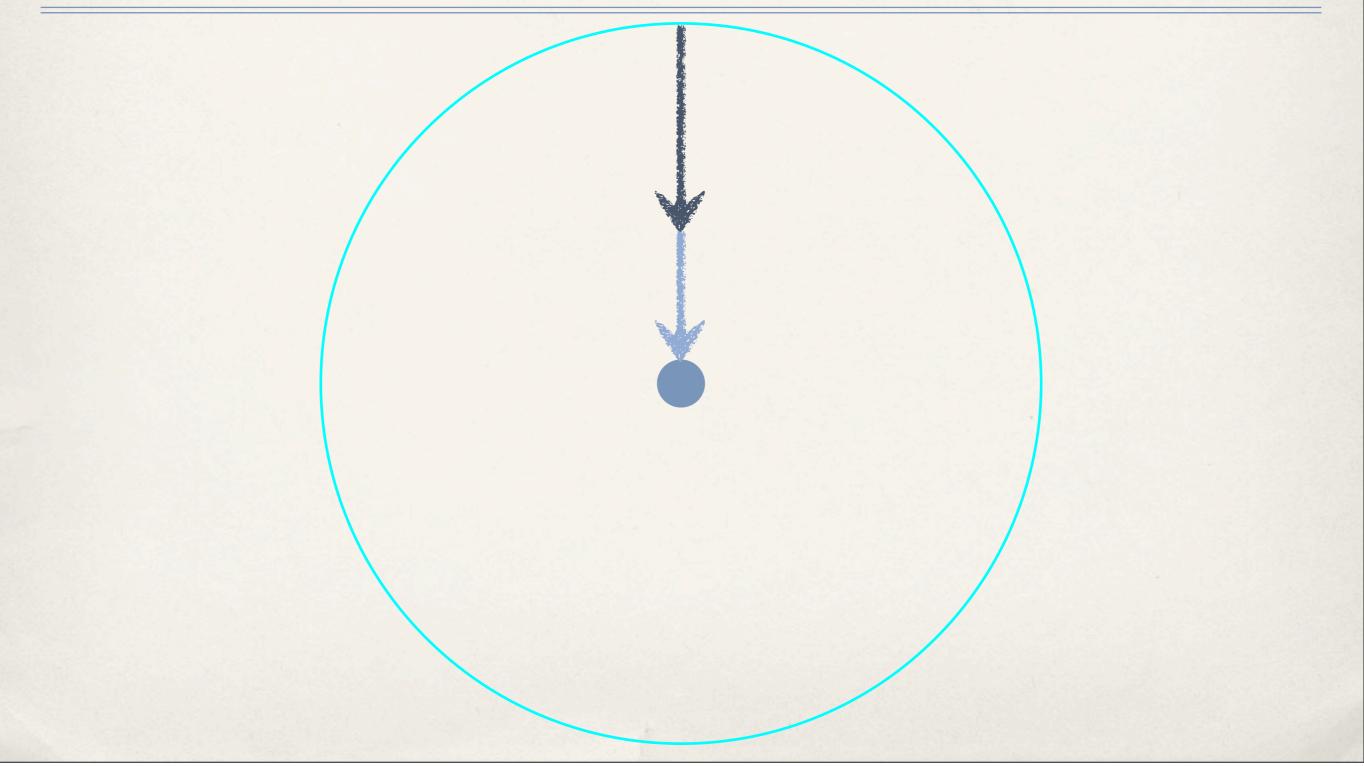


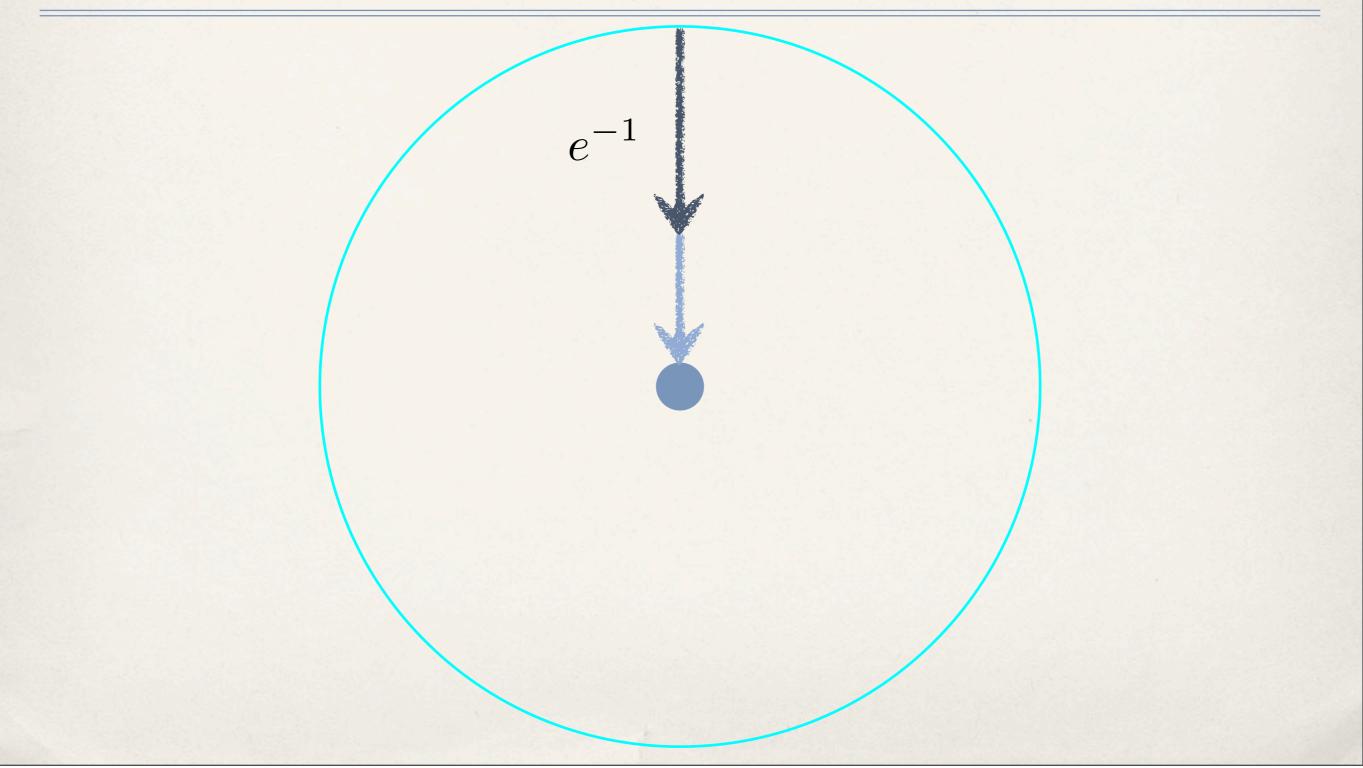


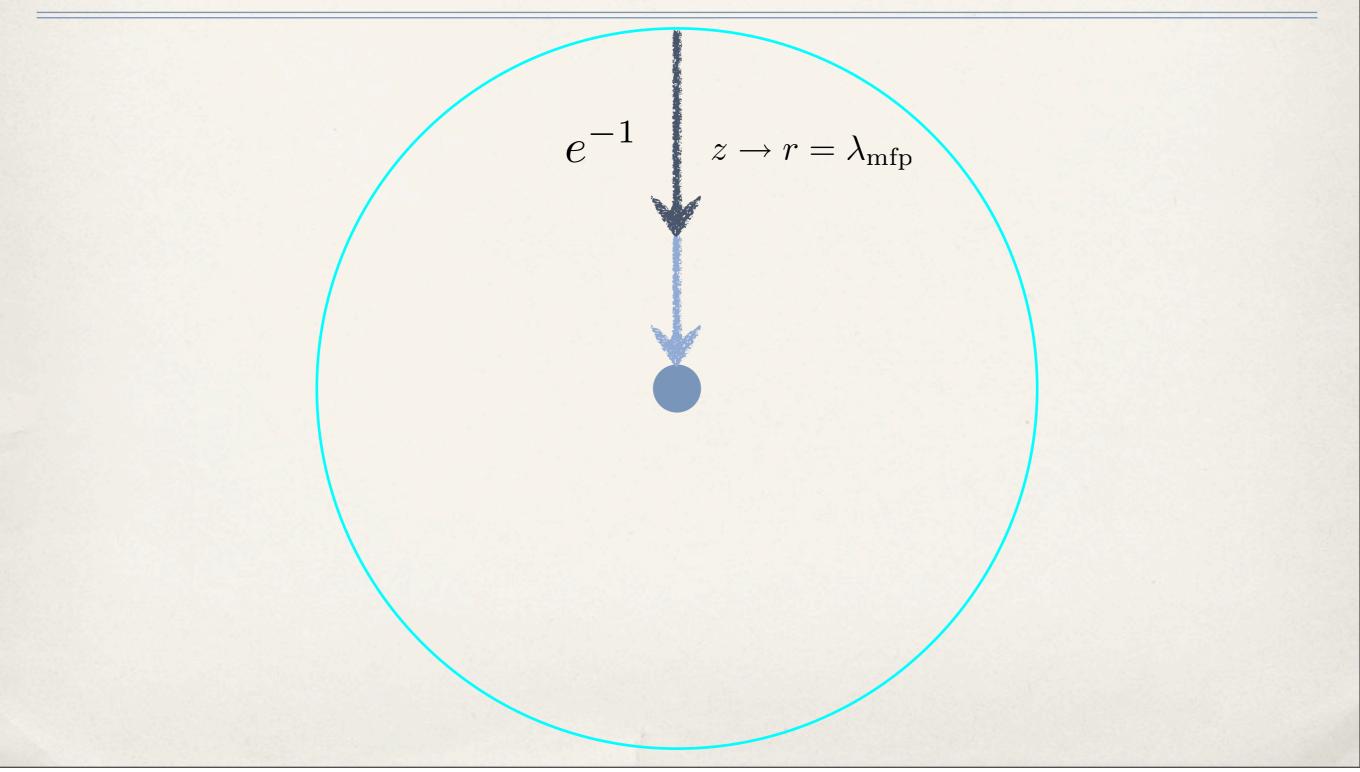


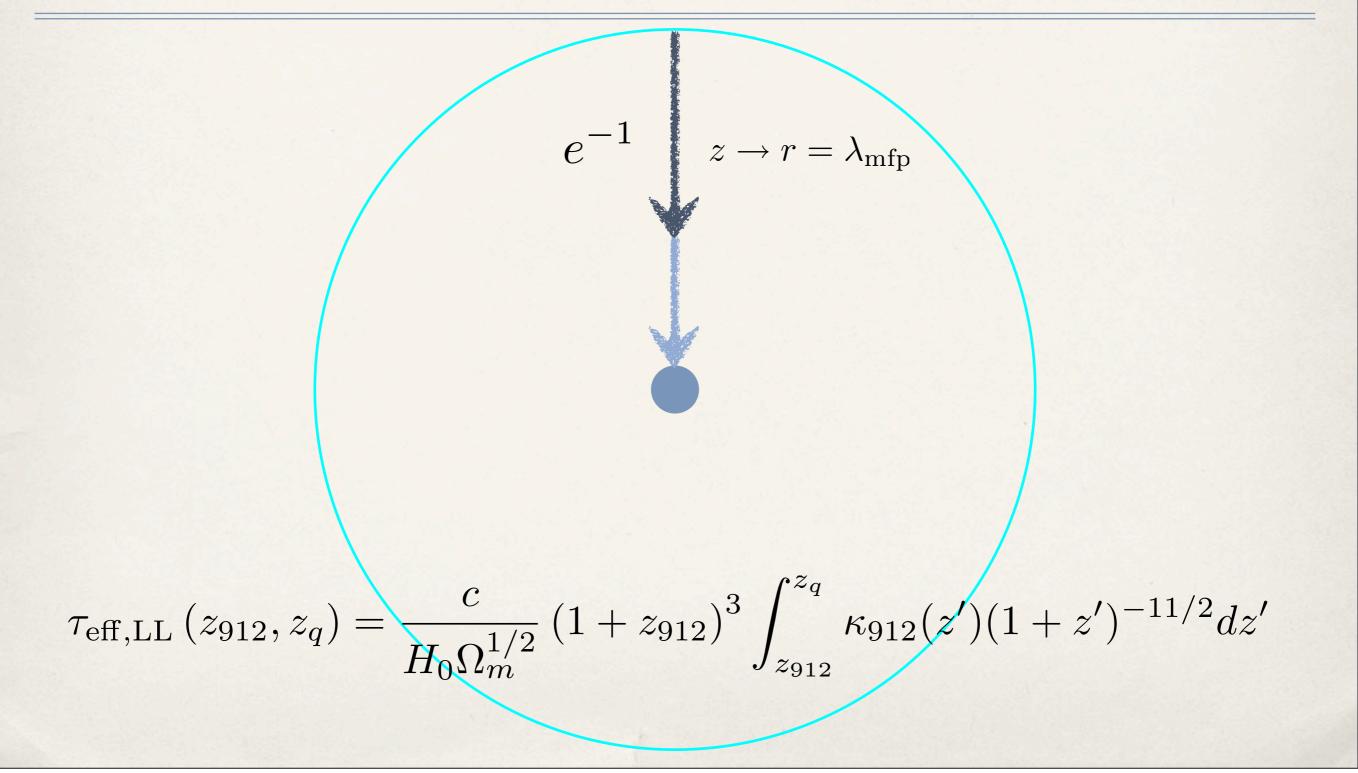




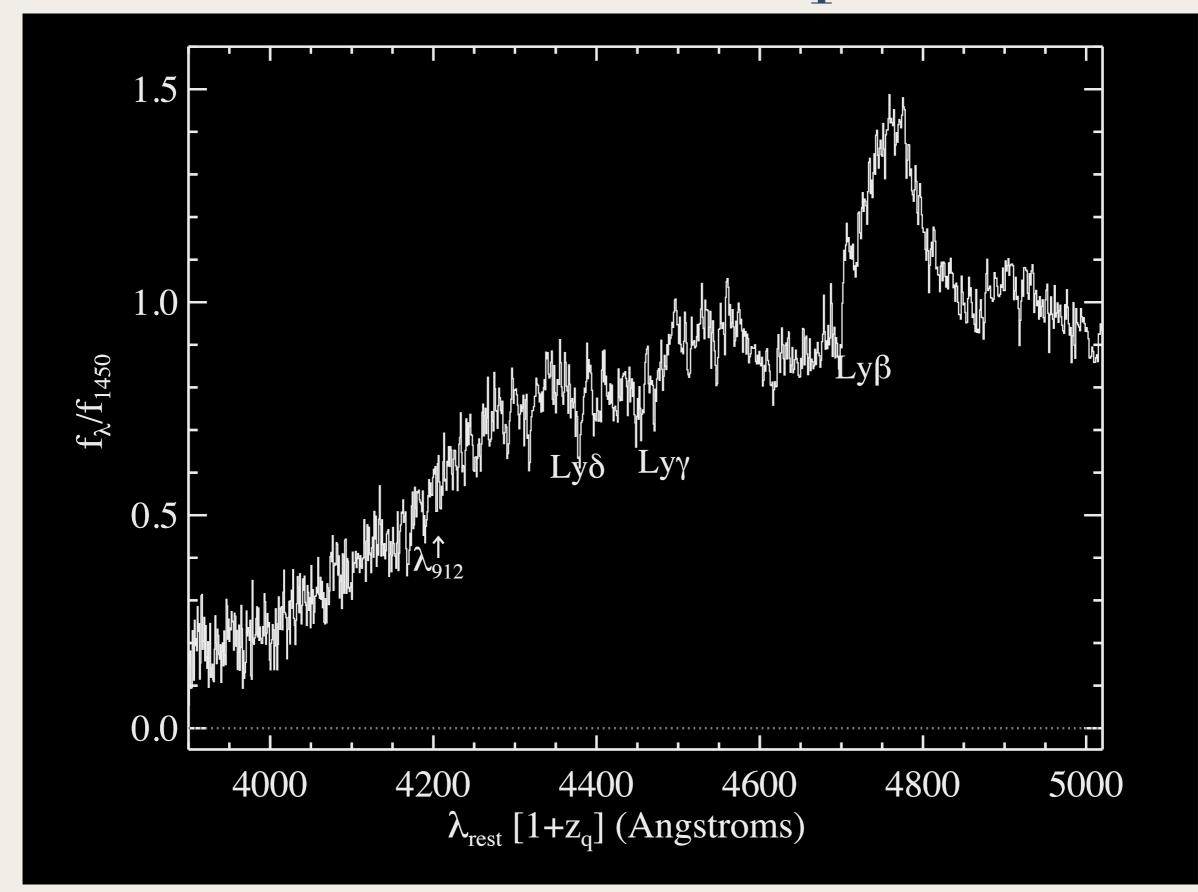




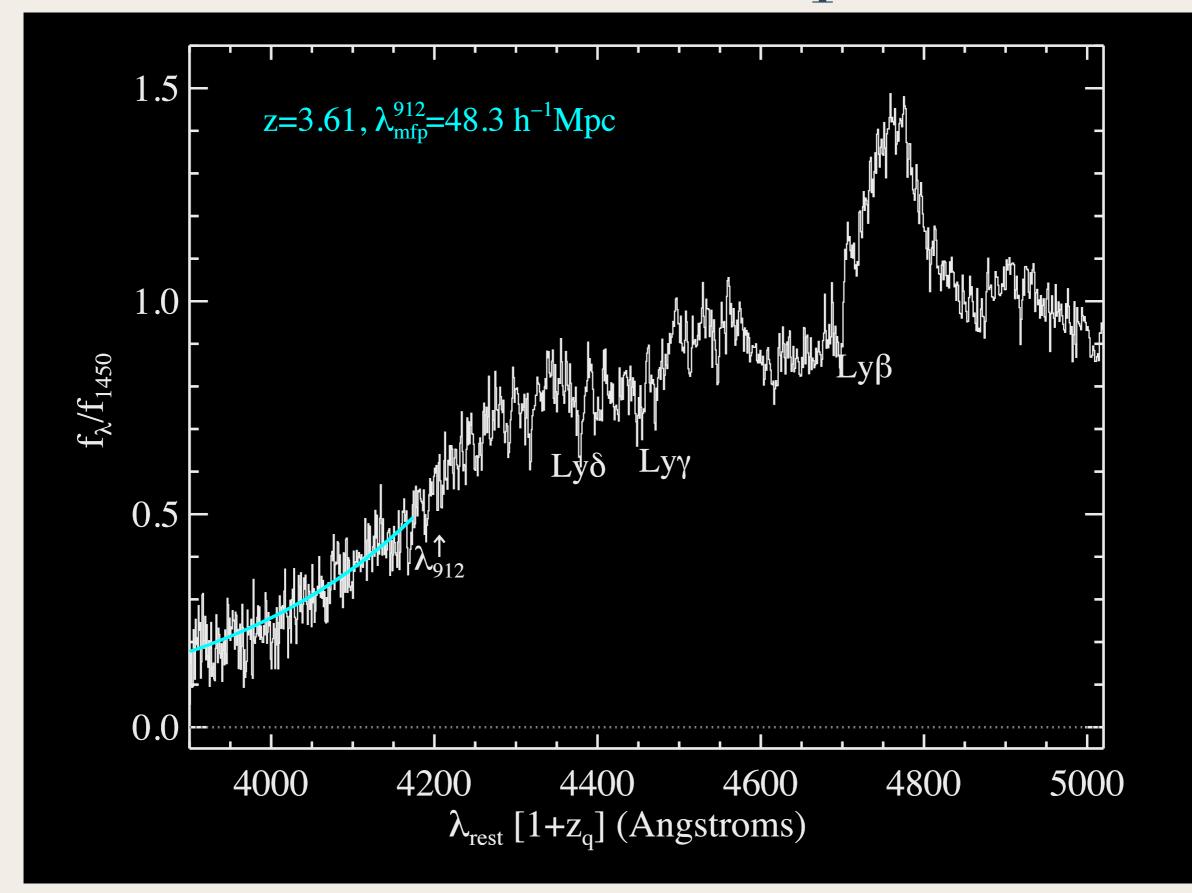




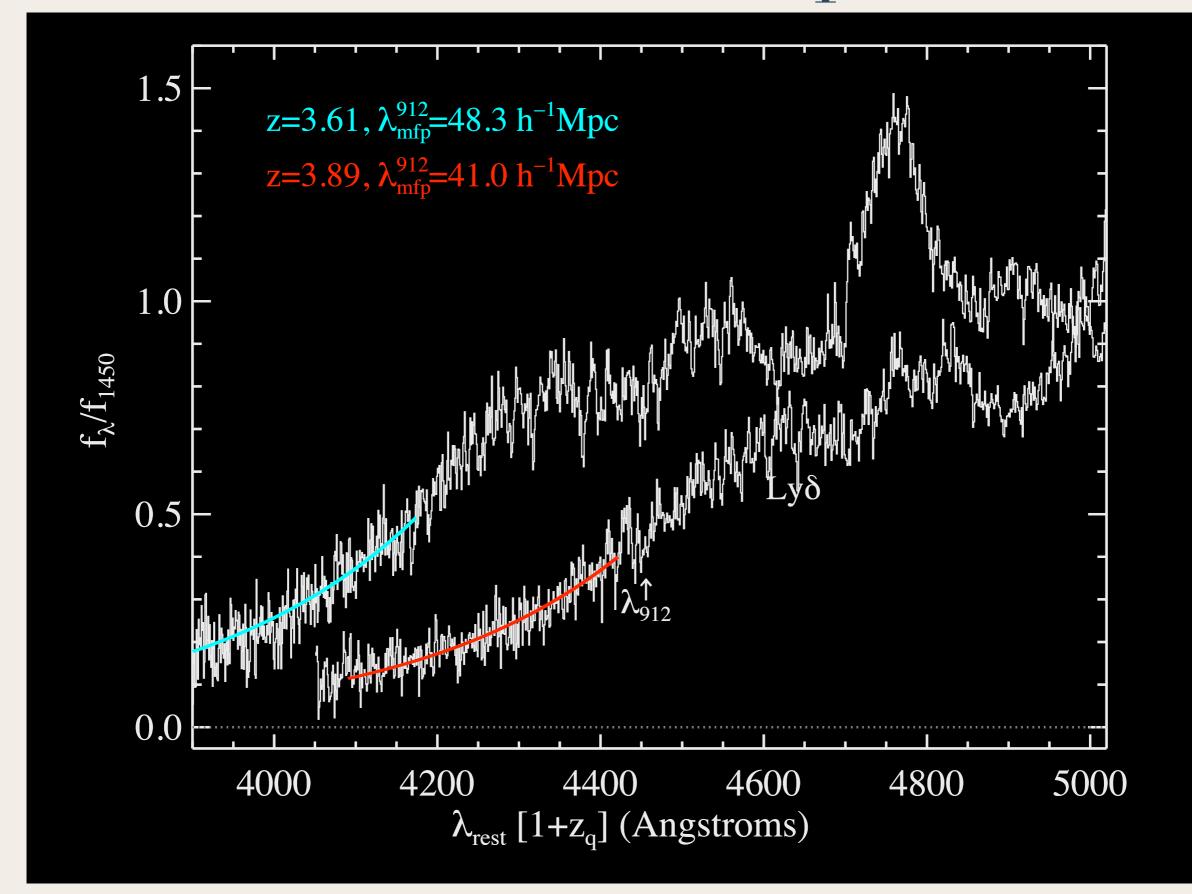
SDSS Stacked Spectra



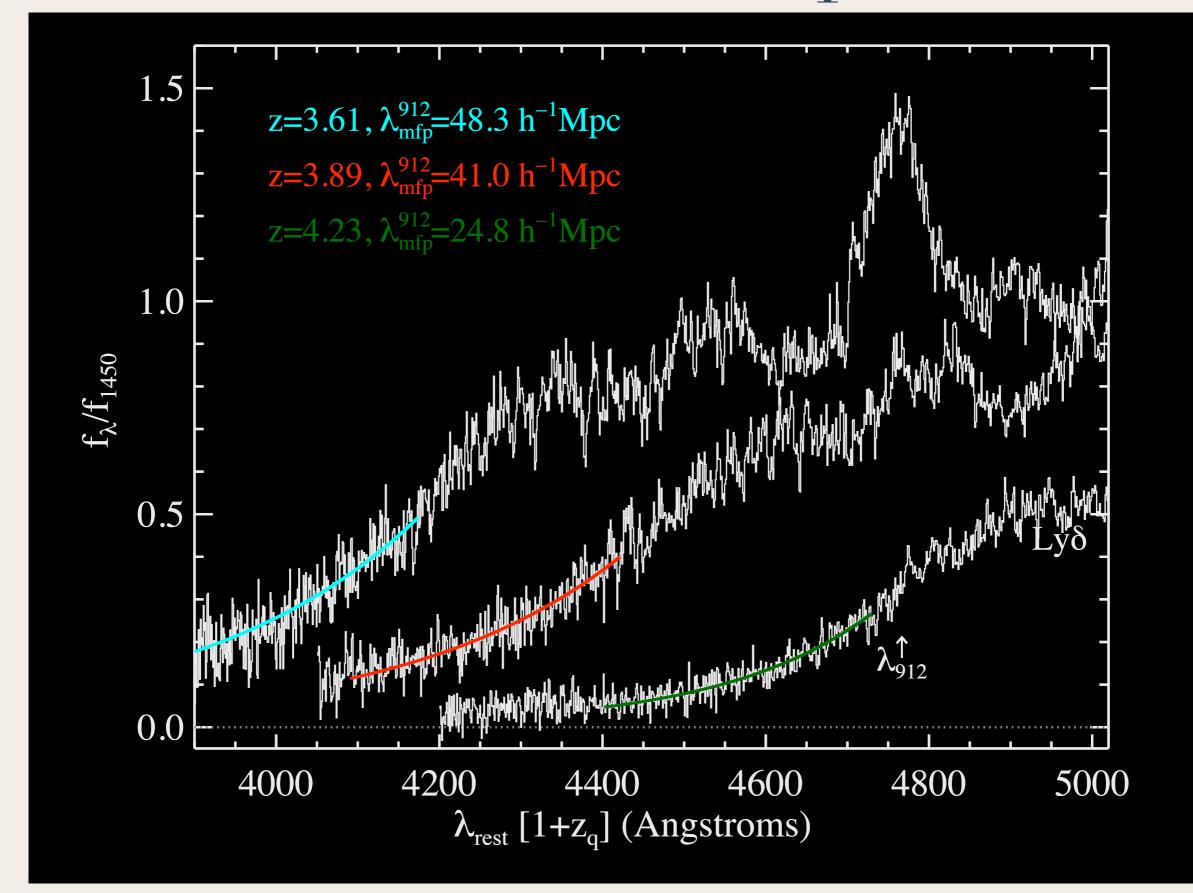
SDSS Stacked Spectra



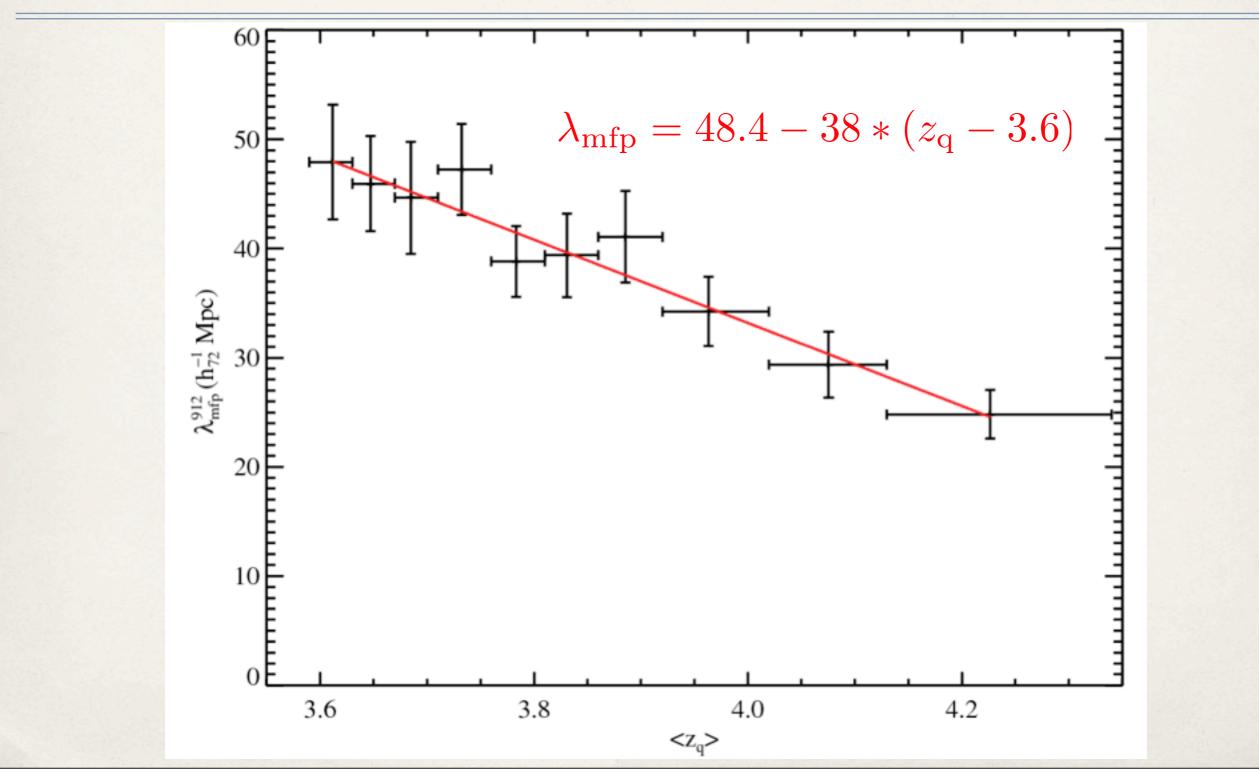
SDSS Stacked Spectra



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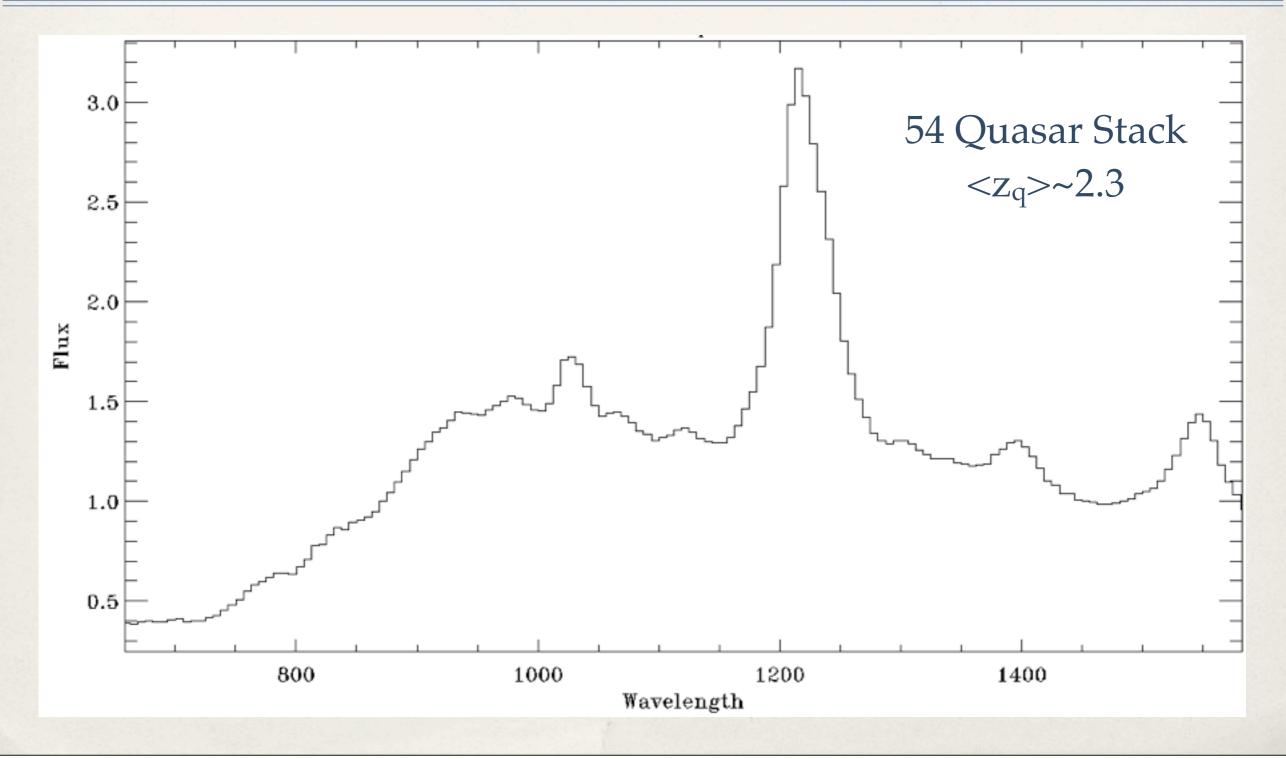
SDSS MFP



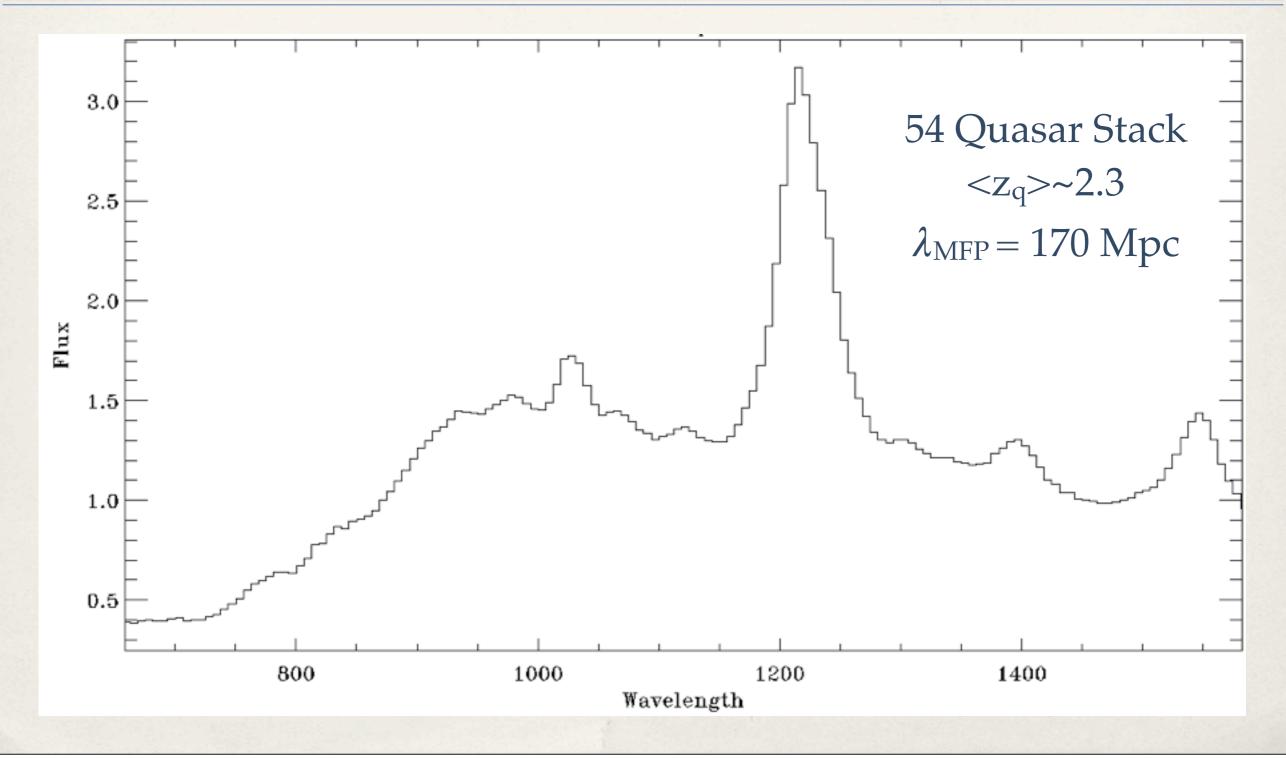
WFC3 Stack

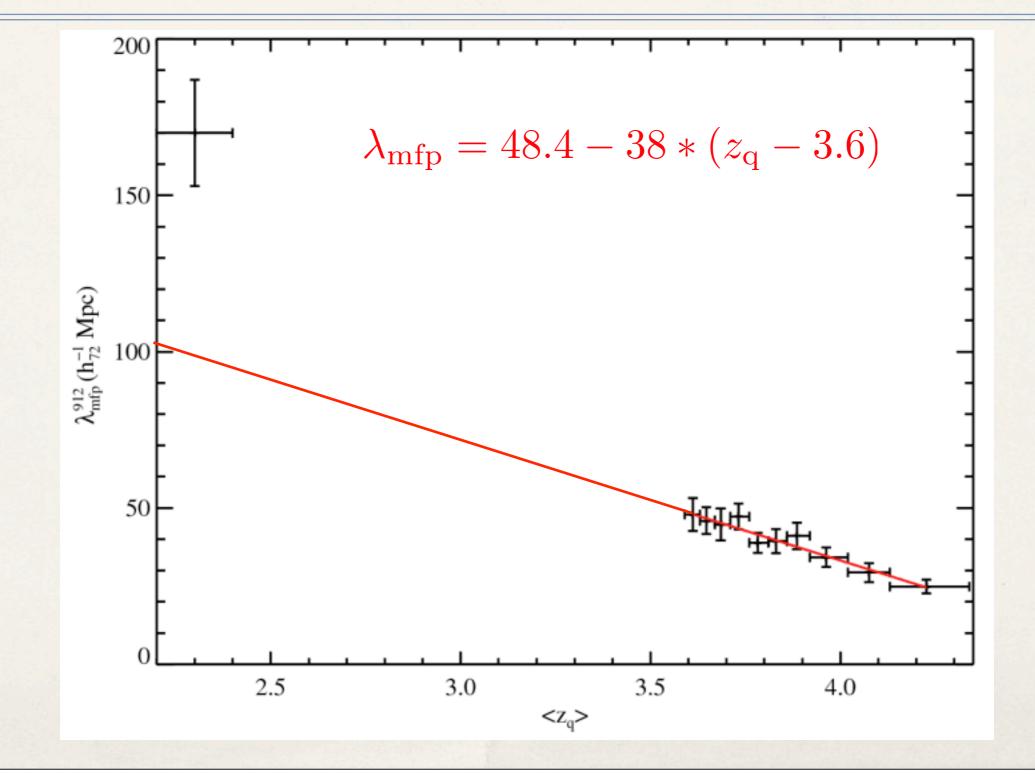
* Can we play the same game at lower z?

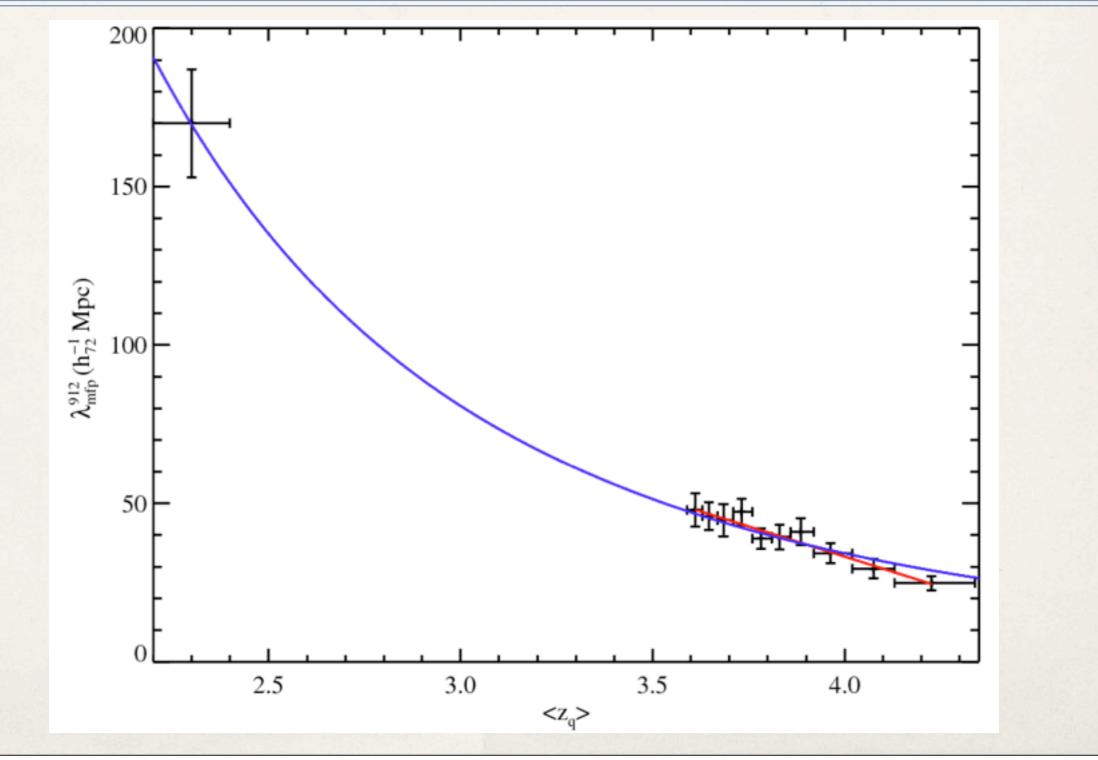
OMG w00t!

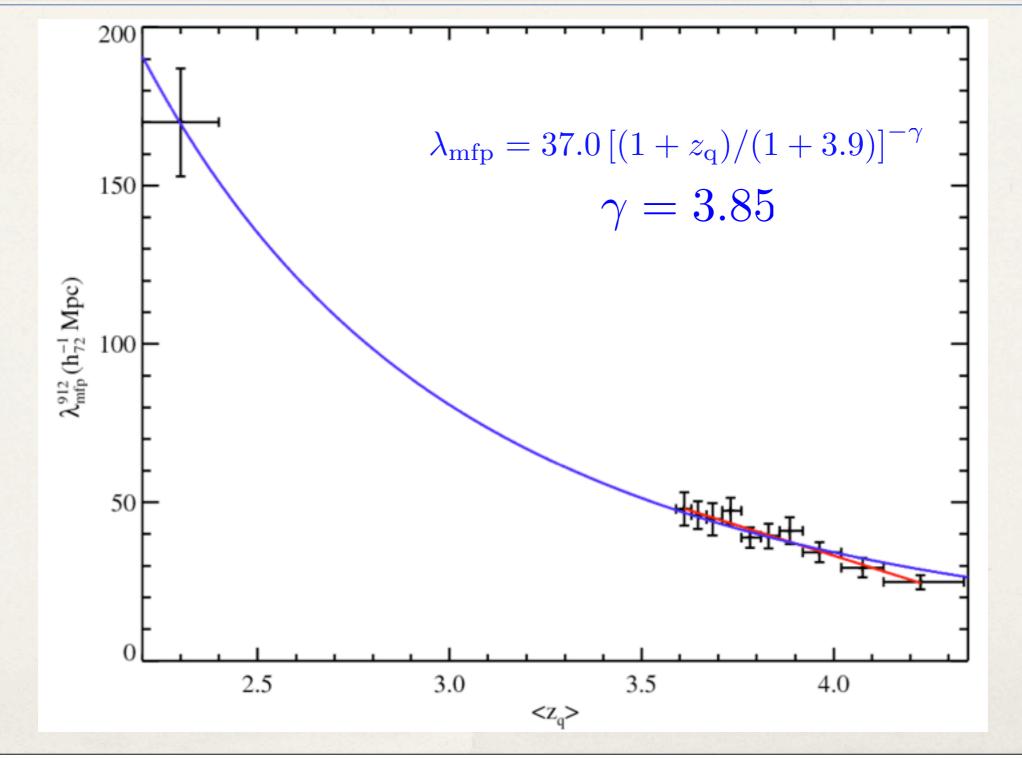


OMG w00t!









* Refine the l(z) analysis, incorporating ACS data

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- * f(N,z) for the partial LLS

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- * f(N,z) for the partial LLS
- * Do the stack right!

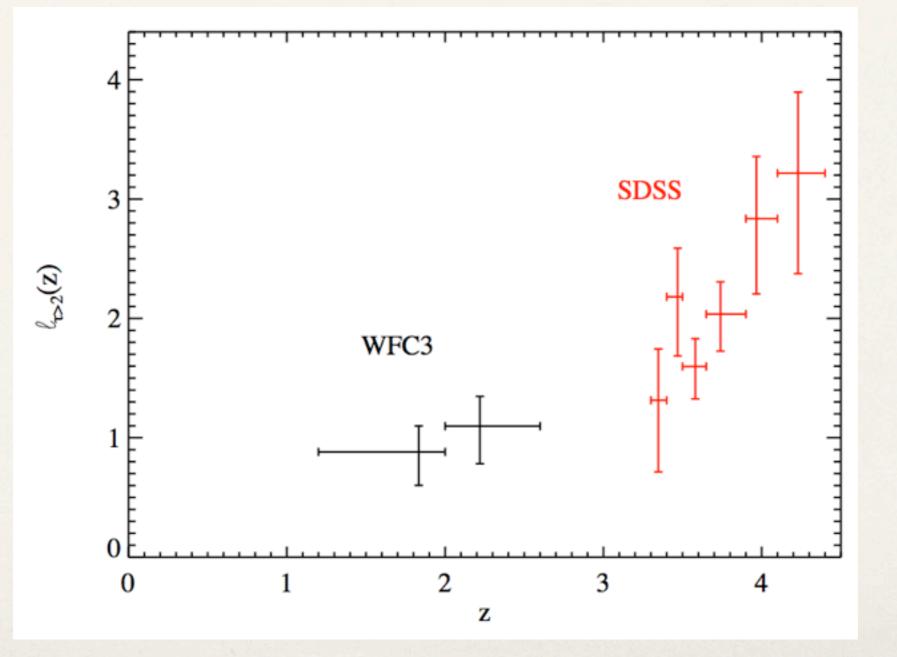
Questions

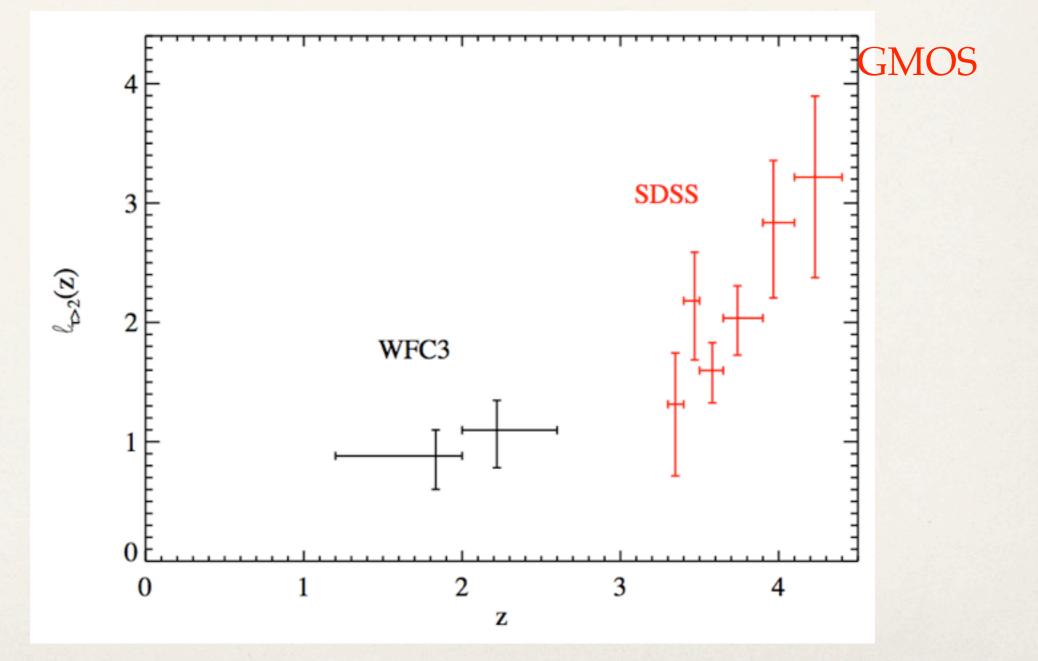
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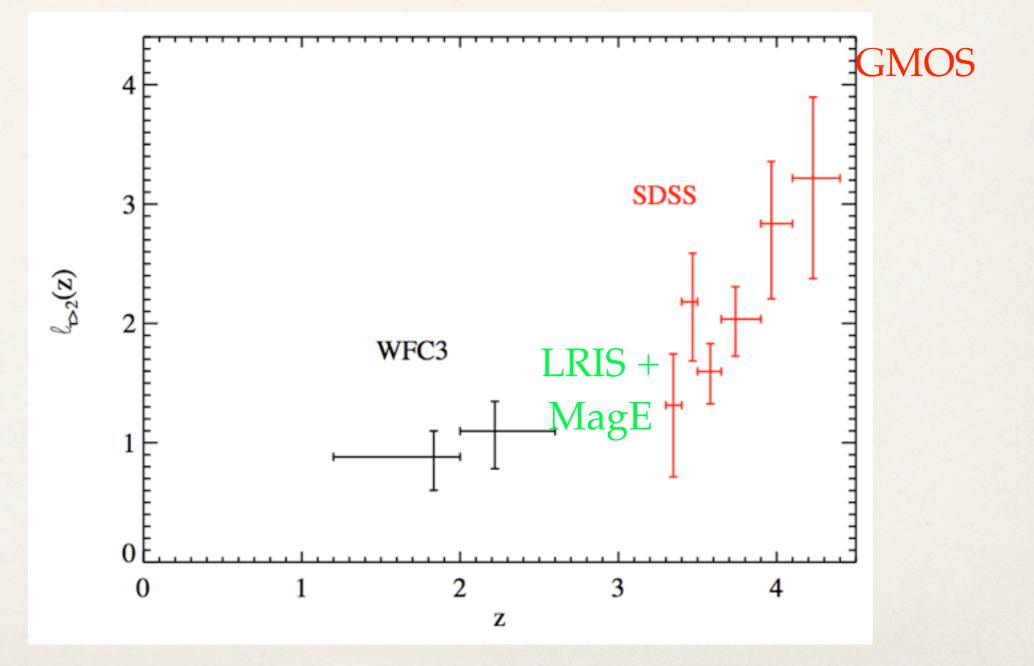
* How does the MFP evolve in z?

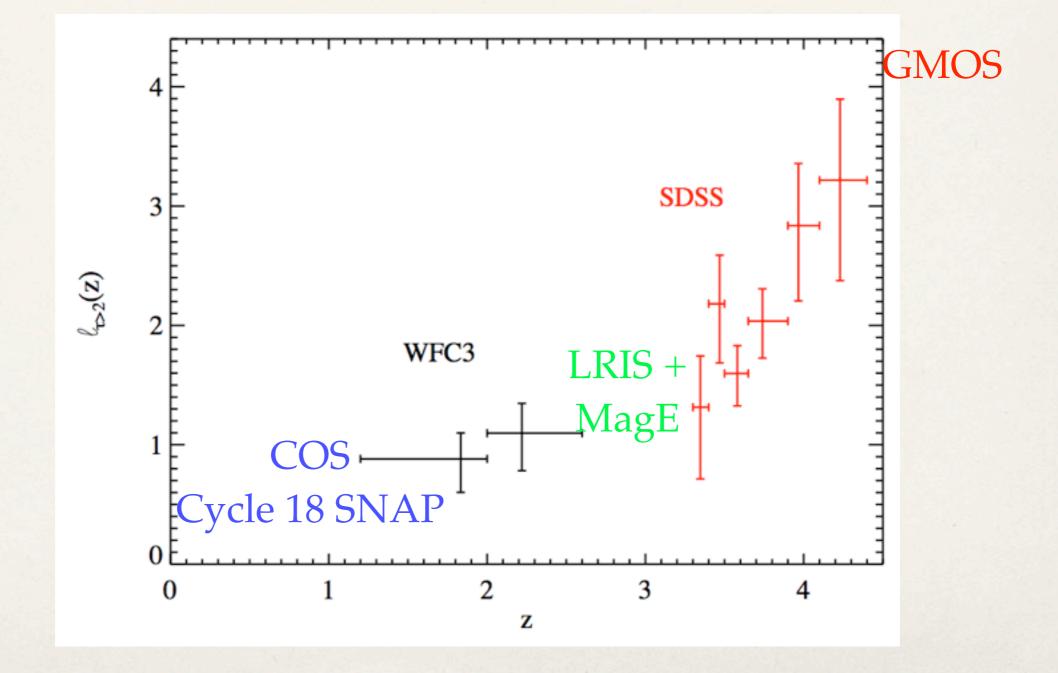
Questions

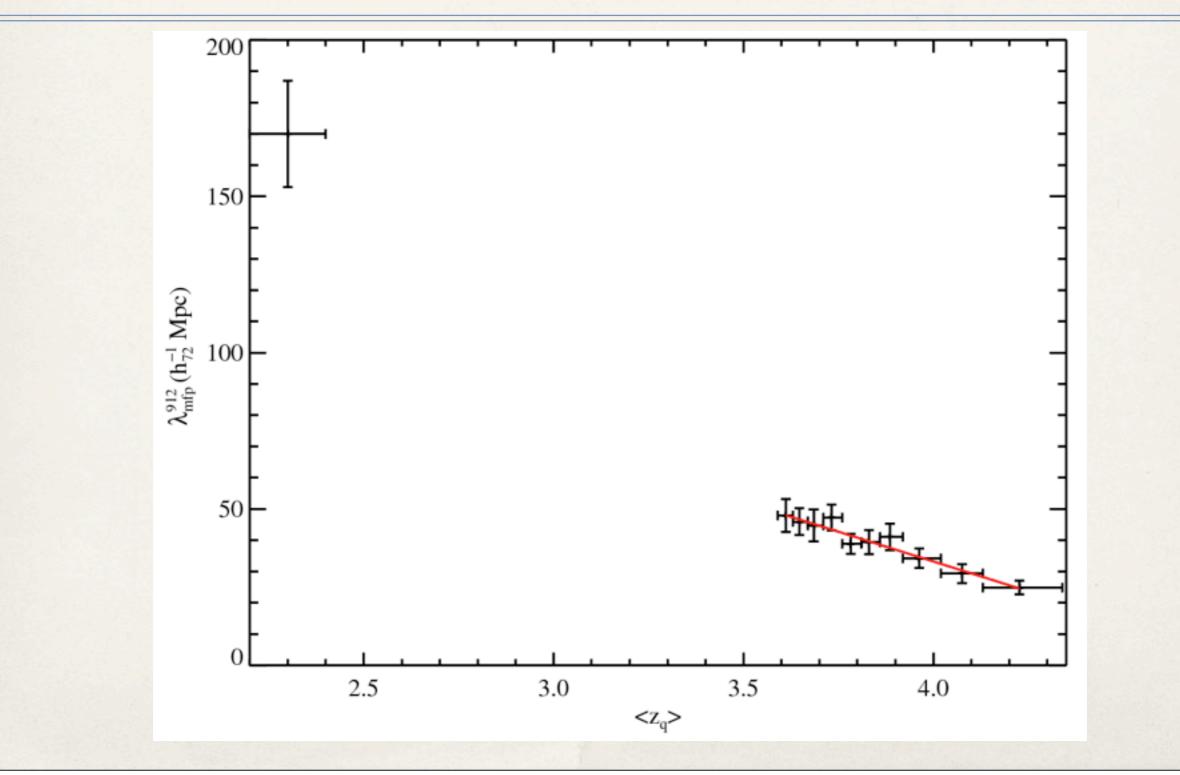
- * How does the MFP evolve in z?
- * How do l(z) and f(N,z) evolve in z?

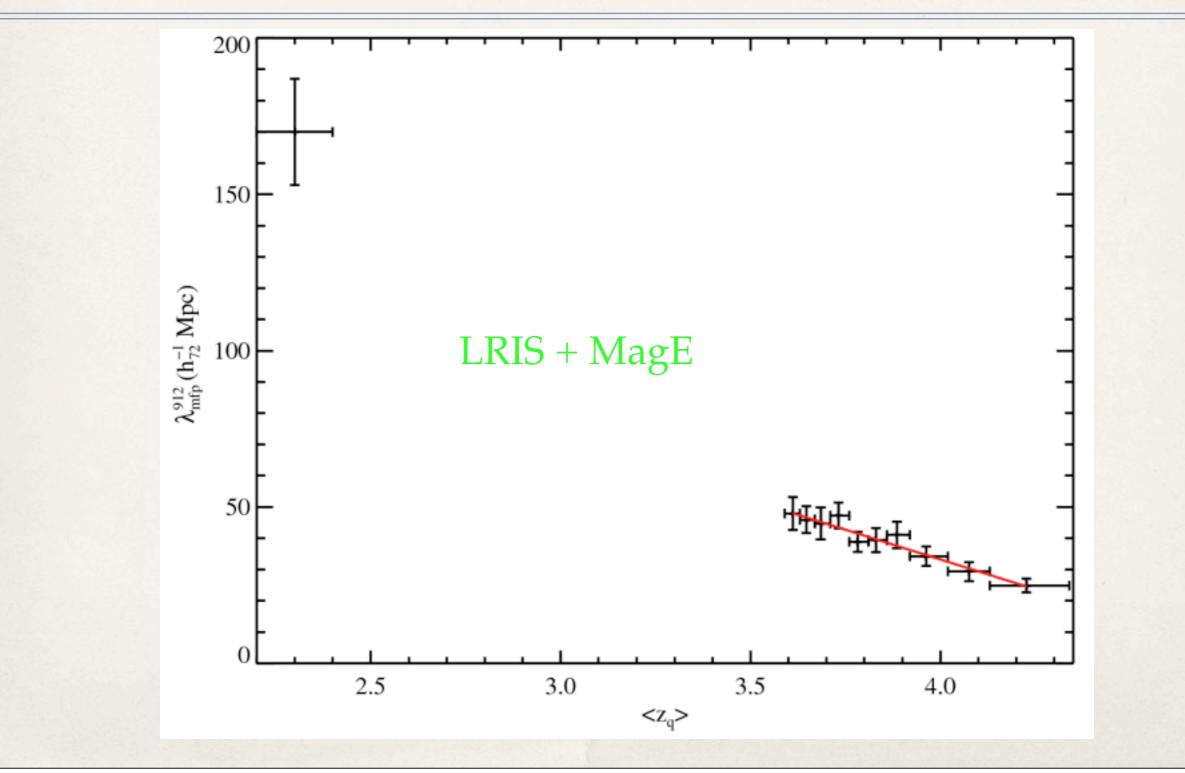


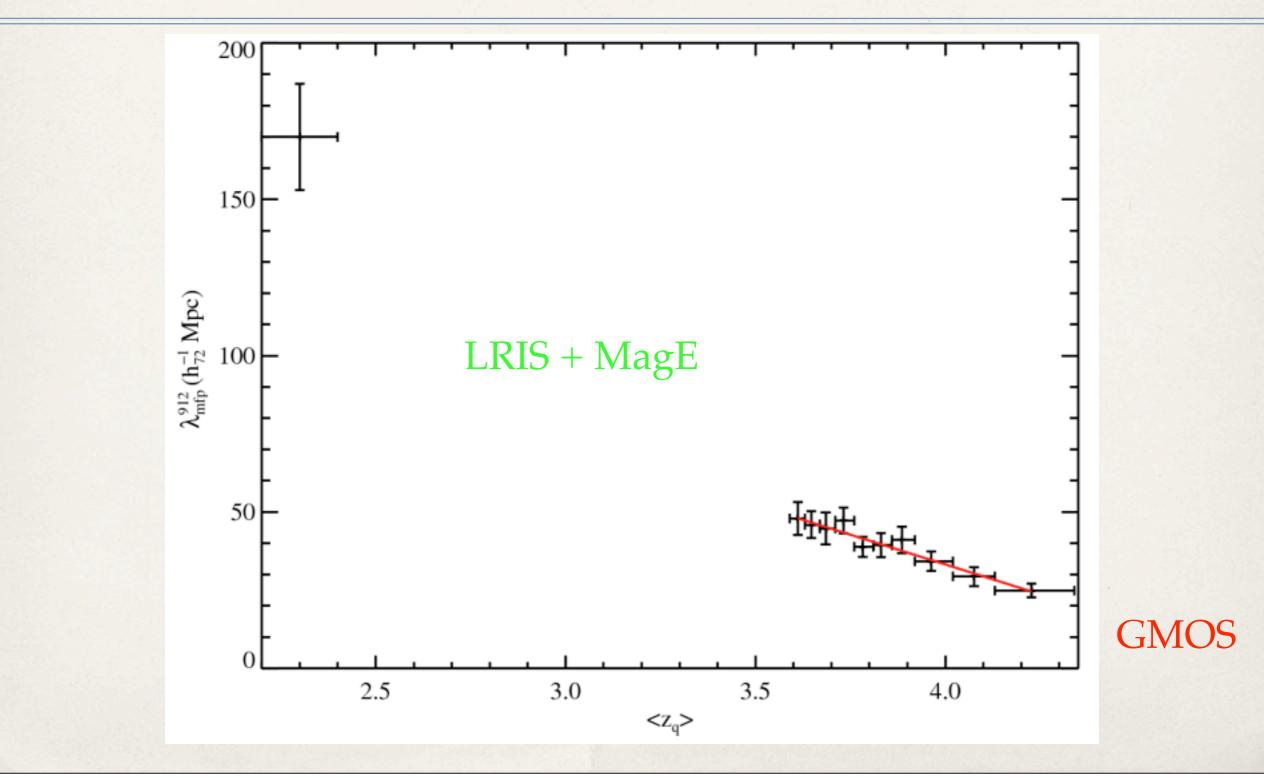




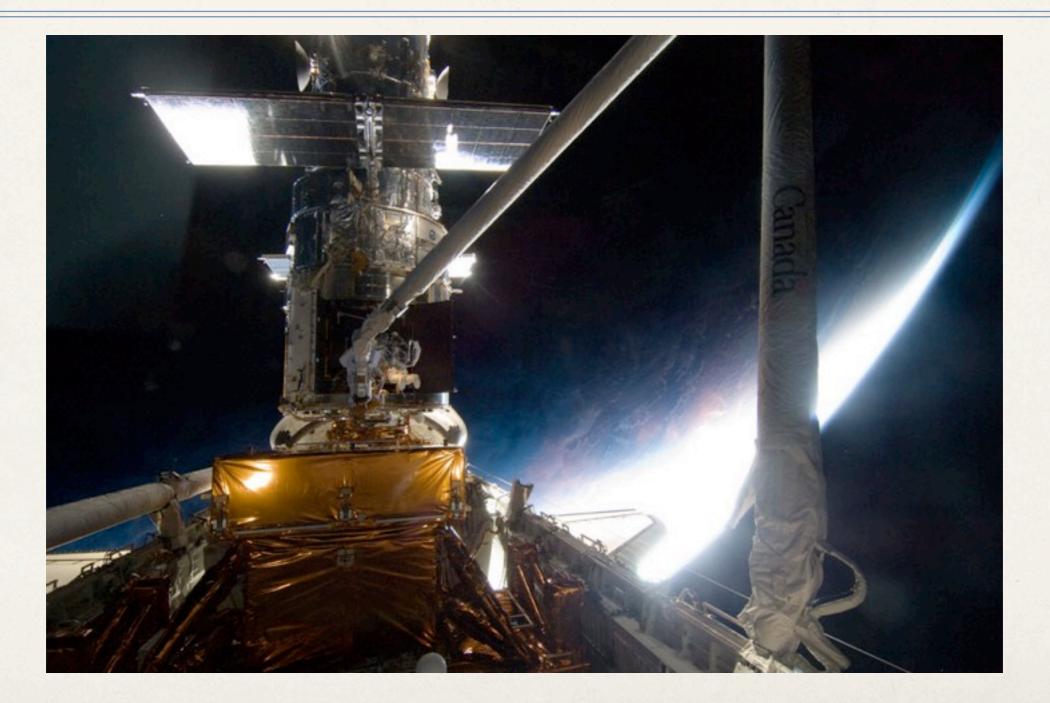




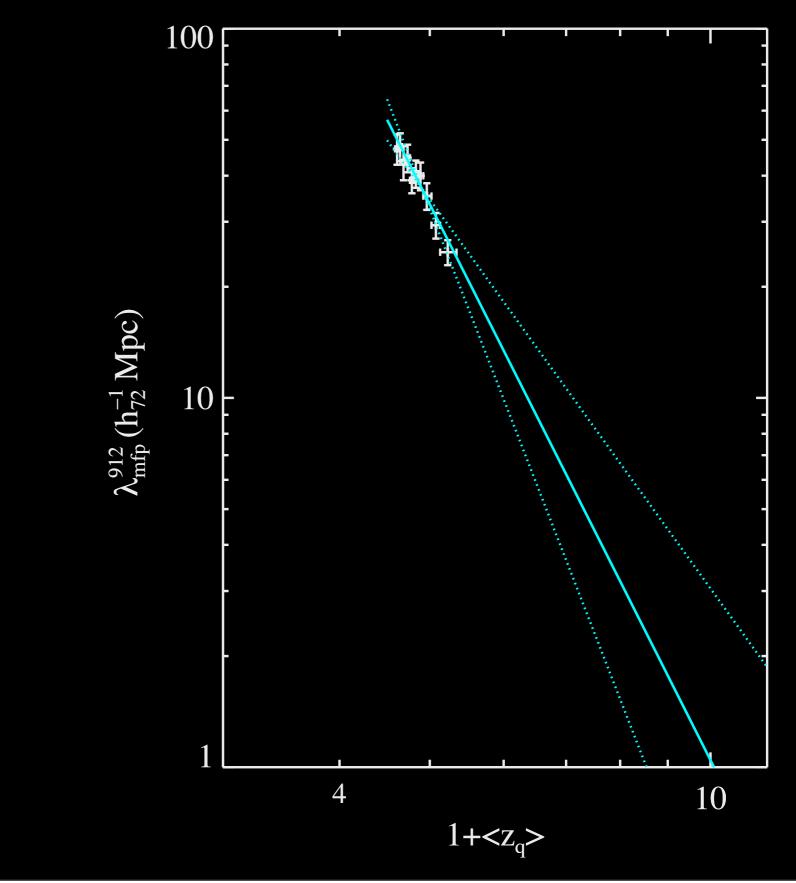




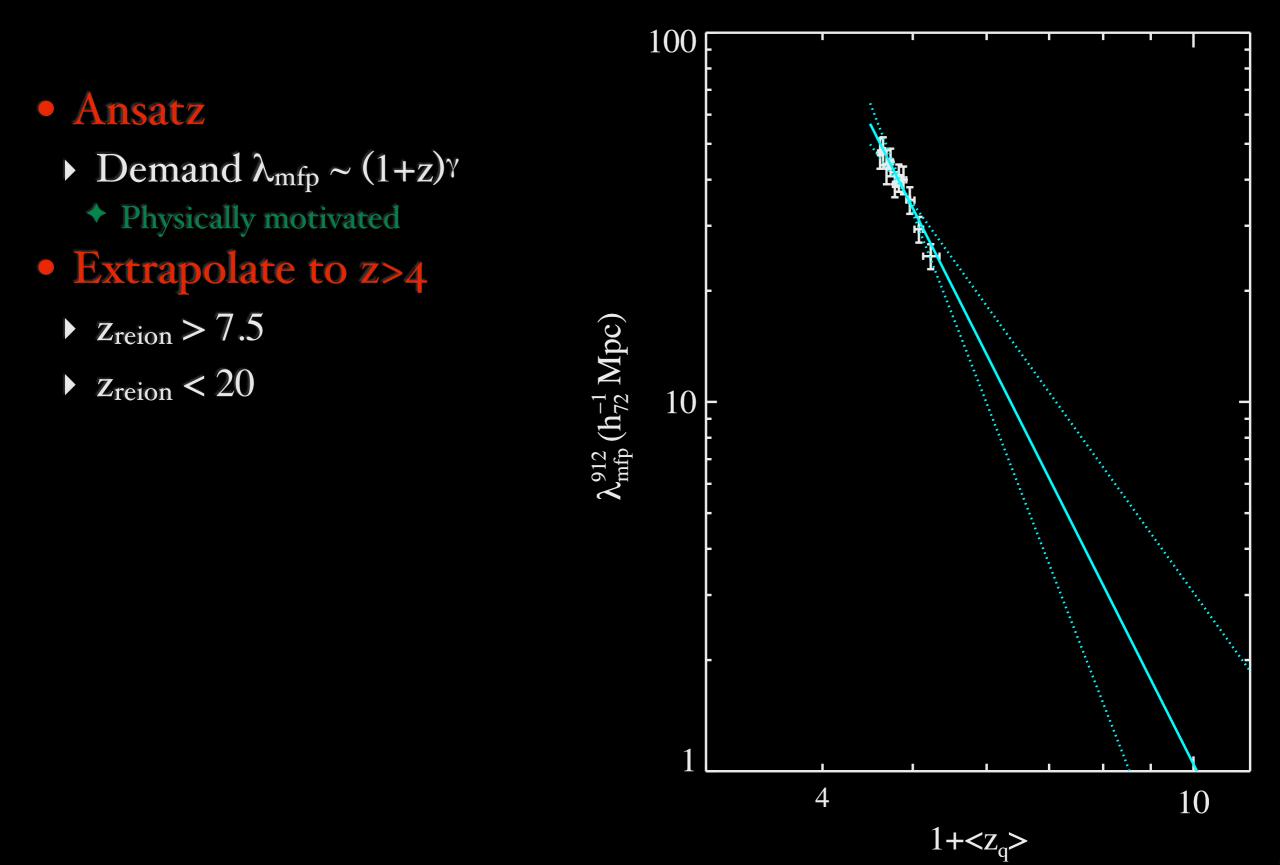
Thanks!



λ_{mfp} Implications: Reionization



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